

Chaos and symmetry breaking in low-dimensional AdS/CFT

Bay Area Particle Theory Seminar - March 2017

Kristan Jensen (SFSU)

based on:

- [arXiv:1605.06098](https://arxiv.org/abs/1605.06098) (PRL 117 (2016), 11, 116601)



An old hope

Question: what is the minimal example of AdS/CFT?

Equivalently: what is the simplest consistent theory
of quantum gravity on AdS?

Pure 3d gravity

Natural candidate: 3d **pure** gravity with negative cc

$$S_{3d} = \frac{1}{16\pi G} \int d^3x \sqrt{-g} \left(R + \frac{2}{L^2} \right)$$

Classical 3d gravity is simple:

Pure 3d gravity

Natural candidate: 3d **pure** gravity with negative cc

Classical 3d gravity is simple:

1. All solutions classified: boundary gravitons plus BTZ
2. May be recast as $SO(2, 1) \times SO(2, 1)$

Chern-Simons theory

$$c = \frac{3L}{2G} = 24k$$

$$S_{3d} = \frac{k}{4\pi} \int \text{tr} \left(A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right)$$

Pure 3d gravity

Natural candidate: 3d **pure** gravity with negative cc

$$S_{3d} = \frac{k}{4\pi} \int \text{tr} \left(A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right) \quad c = \frac{3L}{2G} = 24k$$

Witten 2007: Conjecture that for $k=1,2,\dots$, dual to

“extremal CFT” with torus $Z = |\chi_k(q)|^2$

$$\chi_k(q) = q^{-k} \prod_{n=2}^{\infty} (1 - q^n)^{-1} + O(q)$$

Pure 3d gravity

Natural candidate: 3d **pure** gravity with negative cc

$$Z = |\chi_k(q)|^2 \quad c = \frac{3L}{2G} = 24k$$

k=1: $\chi_1(q) = J(q)$ Modular J-function

this Z corresponds to known CFT! [Frenkel, Lepowsky, Meurman]

So-called Monster CFT

Pure 3d gravity

Natural candidate: 3d **pure** gravity with negative cc

$$Z = |\chi_k(q)|^2 \quad c = \frac{3L}{2G} = 24k$$

k=1: $\chi_1(q) = J(q)$ Modular J-function

Lightest states: 196883 operators of dimension 2

Identifying as BH creation ops, get

$$S = \ln 196883 \approx 12.19 \quad \text{vs.} \quad S_{BH}(k=1) = 4\pi \approx 12.57$$

Pure 3d gravity

Natural candidate: 3d **pure** gravity with negative cc

$$Z = |\chi_k(q)|^2 \quad c = \frac{3L}{2G} = 24k$$

$k=1$: $\chi_1(q) = J(q)$ Modular J-function

$k>1$: not known if corresponding CFT exists,
some evidence against

What about $\text{AdS}_2/\text{CFT}_1$?

Simplest possible setting for the duality

And ubiquitous: AdS_2 generically appears as
near-horizon of SUSY black holes

What about $\text{AdS}_2/\text{CFT}_1$?

Simplest possible setting for the duality

And ubiquitous: AdS_2 generically appears as
near-horizon of SUSY black holes

PROBLEM: Neither AdS_2 gravity nor CFT_1 exist

The problem

AdS₂: cannot support finite-energy excitations, $\langle E \rangle = 0$

CFT₁: One runs into a paradox of [Polchinski]

scale-invariant density of states in 1d is

$$\rho(E) = e^{S_1} \delta(E) + \frac{e^{S_2}}{E}$$

The problem

scale-invariant density of states in 1d is

$$\rho(E) = \boxed{e^{S_1} \delta(E)} + \frac{e^{S_2}}{E}$$




Zero-energy states
No dynamics!

The problem

scale-invariant density of states in 1d is

$$\rho(E) = e^{S_1} \delta(E) + \frac{e^{S_2}}{E}$$



Divergent at E=0
Z does not exist!

Goal for this talk:

Obtain sensible notion of $\text{AdS}_2/\text{CFT}_1$

Outline

1. Motivation
2. $\text{NAdS}_2/\text{NCFT}_1$
3. Chaos
4. Parting words

Outline

1. Motivation
- 2. $\text{NAdS}_2/\text{NCFT}_1$**
3. Chaos
4. Parting words

Dilaton gravity

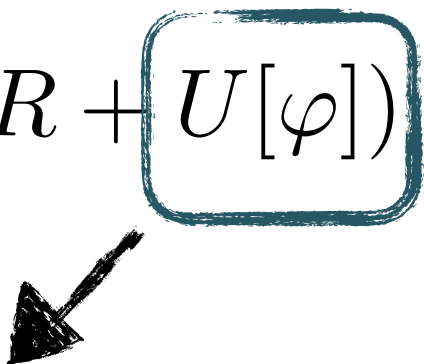
Pure 2d gravity is topological $\int d^2x \sqrt{-g} R = 4\pi\chi$

Reduction to 2d leads to **dilaton** gravity

$$S = \frac{1}{2\kappa^2} \int d^2x \sqrt{-g} (\varphi R + U[\varphi]) + S_{\text{matter}}$$

Dilaton gravity

Reduction to 2d leads to dilaton gravity

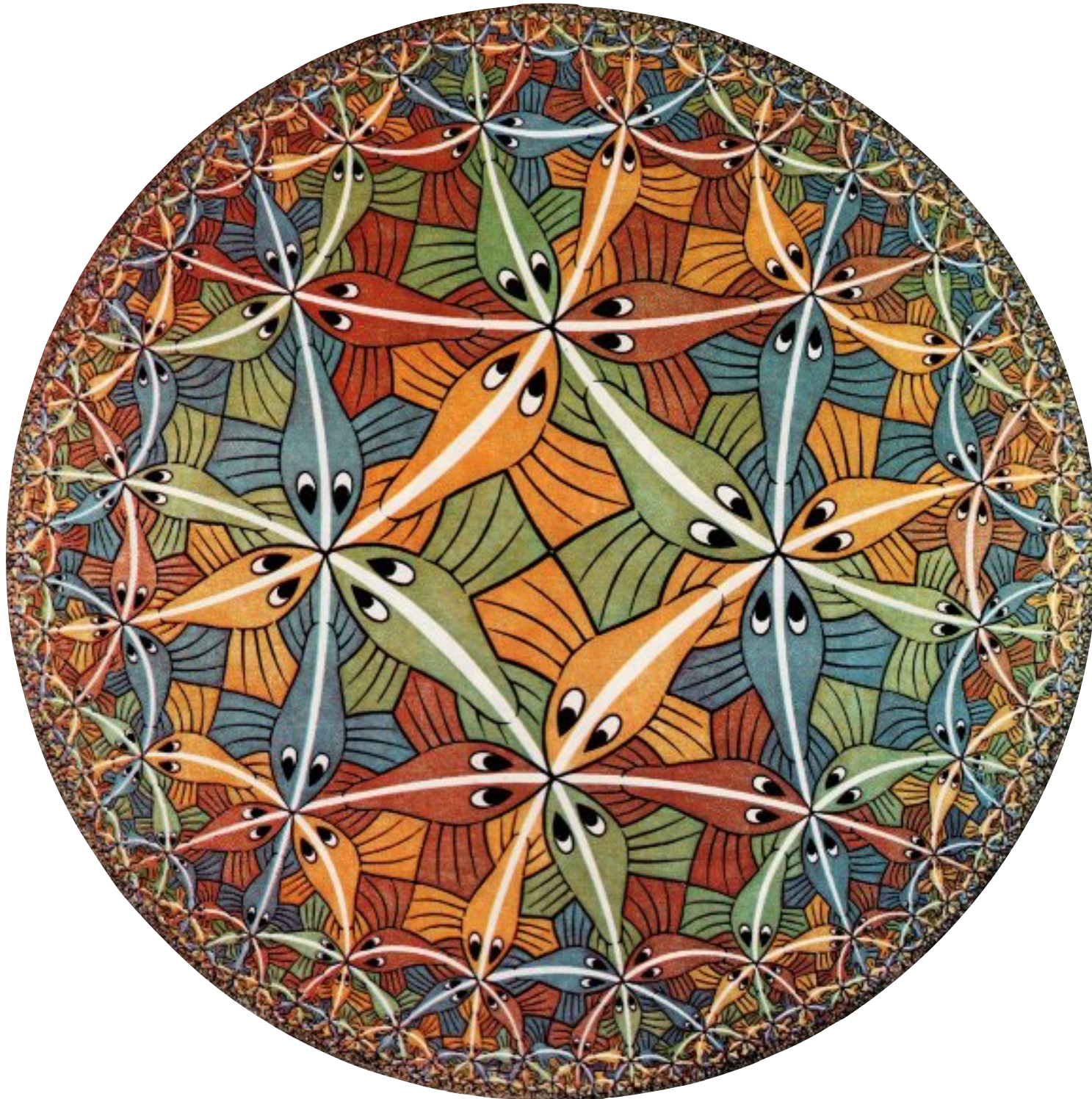
$$S = \frac{1}{2\kappa^2} \int d^2x \sqrt{-g} (\varphi R + U[\varphi]) + S_{\text{matter}}$$


General two-derivative theory
characterized by U

AdS₂ solutions with constant dilaton at roots of U

$$U[\varphi_0] = 0, \quad U'[\varphi_0] = \frac{2}{L^2}$$

AdS₂



Moduli of AdS_2

Most general AdS_2 spacetime characterized by a free function of boundary time

$$g = -r^2 \left(1 + \frac{h(t)}{r^2} \right)^2 dt^2 + \frac{dr^2}{r^2} \quad \varphi = \varphi_0$$

Moduli of AdS_2


Most general AdS_2 spacetime characterized by a free function of boundary time

Convenient to redefine $h(t) = \frac{1}{2}\{f(t), t\}$

Moduli of AdS_2

Most general AdS_2 spacetime characterized by a free function of boundary time

Convenient to redefine $h(t) = \frac{1}{2} \{f(t), t\}$

- 
1. $f(t)$ acts as conformal transformation on bdy
 2. For any $f(t)$, invariant under $PSL(2, \mathbb{R})$

$$f(t) \rightarrow \frac{af(t) + b}{cf(\tau) + d}$$

Moduli of AdS_2

Convenient to redefine $h(t) = \frac{1}{2}\{f(t), t\}$

1. $f(t)$ acts as conformal transformation on bdy
2. For any $f(t)$, invariant under $PSL(2, \mathbb{R})$

boundary conformal symmetry
= Virasoro symmetry with $c=0$

NAdS₂ spacetimes

Name due to [Maldacena, Stanford]

$U[\varphi] \approx 2(\varphi - \varphi_0)$ also admits linear dilaton solutions

NAdS₂ spacetimes

Name due to [Maldacena, Stanford]

$$U[\varphi] \approx 2(\varphi - \varphi_0) \longrightarrow \begin{aligned} g &= -r^2 dt^2 + \frac{dr^2}{r^2} \\ \varphi &= \ell r + \varphi_0 \end{aligned}$$

Arises as IR endpoint of ANY
holographic RG flow ending in AdS₂

NAdS₂ spacetimes

Name due to [Maldacena, Stanford]

$$U[\varphi] \approx 2(\varphi - \varphi_0) \longrightarrow \begin{aligned} g &= -r^2 dt^2 + \frac{dr^2}{r^2} \\ \varphi &= \ell r + \varphi_0 \end{aligned}$$

Moduli $f(t)$ become pseudo-moduli: Virasoro $\longrightarrow PSL(2, \mathbb{R})$

Dual to nearly conformal QM, dubbed NCFT₁

Fluid/gravity

Write NAdS₂ in infalling Eddington-Finkelstein coordinates

$$g = - \left(r^2 + 2\{f(t), t\} \right) dt^2 + 2dt dr + O(\varepsilon)$$

$$\varphi = \varphi_0 + \varepsilon \ell r + O(\varepsilon^2) \qquad T_{\mu\nu} = O(\varepsilon)$$

Solve bulk eoms exactly near AdS₂; can rewrite as equation on the boundary:

$$\dot{E} = \dot{\lambda} \langle O \rangle$$

with

$$E = -\frac{\ell}{\kappa^2} \{f(t), t\} + (1 - \Delta) \lambda \langle O \rangle$$

Schwarzian action

Both follow from action where $f(t)$ is fundamental field:

$$S = -\frac{\ell}{\kappa^2} \int dt \{f(t), t\} + W[\lambda(t); f(t)]$$

with $\{f(t), t\} = \frac{f'''(t)}{f'(t)} - \frac{3}{2} \left(\frac{f''(t)}{f'(t)} \right)^2 = \text{Schwarzian derivative}$

[KJ], [Maldacena, Stanford, Yang], [Engelsoy, Martens, Verlinde]

also obtained in the low-energy limit of Sachdev-Ye-Kitaev models [Maldacena, Stanford]

Schwarzian action

Both follow from action where $f(t)$ is fundamental field:

$$S = -\frac{\ell}{\kappa^2} \int dt \{f(t), t\} + W[\lambda(t); f(t)] + O(\ell\lambda^2, \ell^2)$$

Leading interactions consistent with

$$\text{Virasoro} \longrightarrow PSL(2, \mathbb{R})$$

The good and the ugly

(Very) good:

Rewrite two-derivative NAdS₂ gravity as 1d theory

The good and the ugly

(Very) good:

Rewrite two-derivative NAdS₂ gravity as 1d theory

Ugly:

Resulting Schwarzian theory is sick in isolation

[Stanford, Witten]

(situation recalls [Maloney, Witten] prescription for partition function of pure 3d gravity)

Outline

1. Motivation
2. $\text{NAdS}_2/\text{NCFT}_1$
- 3. Chaos**
4. Parting words

Lyapunov exponent

Quantum mechanical analogue of Lyapunov exponent, characterizing early-time chaotic growth:

$$\langle [W(t), V(0)]^2 \rangle_{\beta} \sim e^{\lambda_L t}$$

Lyapunov exponent

Quantum mechanical analogue of Lyapunov exponent, characterizing early-time chaotic growth:

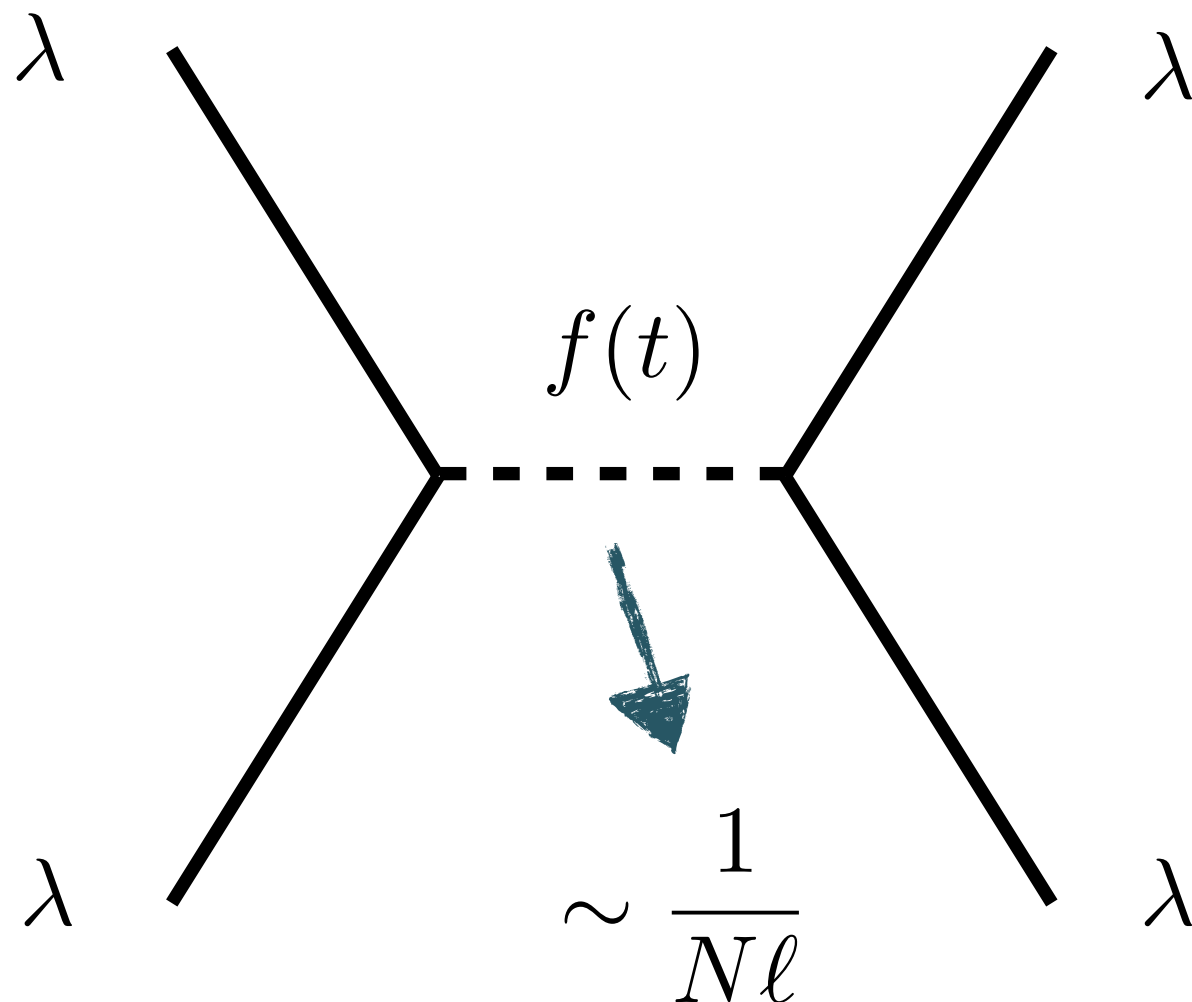
$$\langle [W(t), V(0)]^2 \rangle_\beta \sim e^{\lambda_L t}$$

Two properties:

1. [Maldacena, Shenker, Stanford] $\lambda_L \leq \frac{2\pi}{\beta}$
2. Dual to Einstein gravity is maximally chaotic
[Shenker, Stanford]

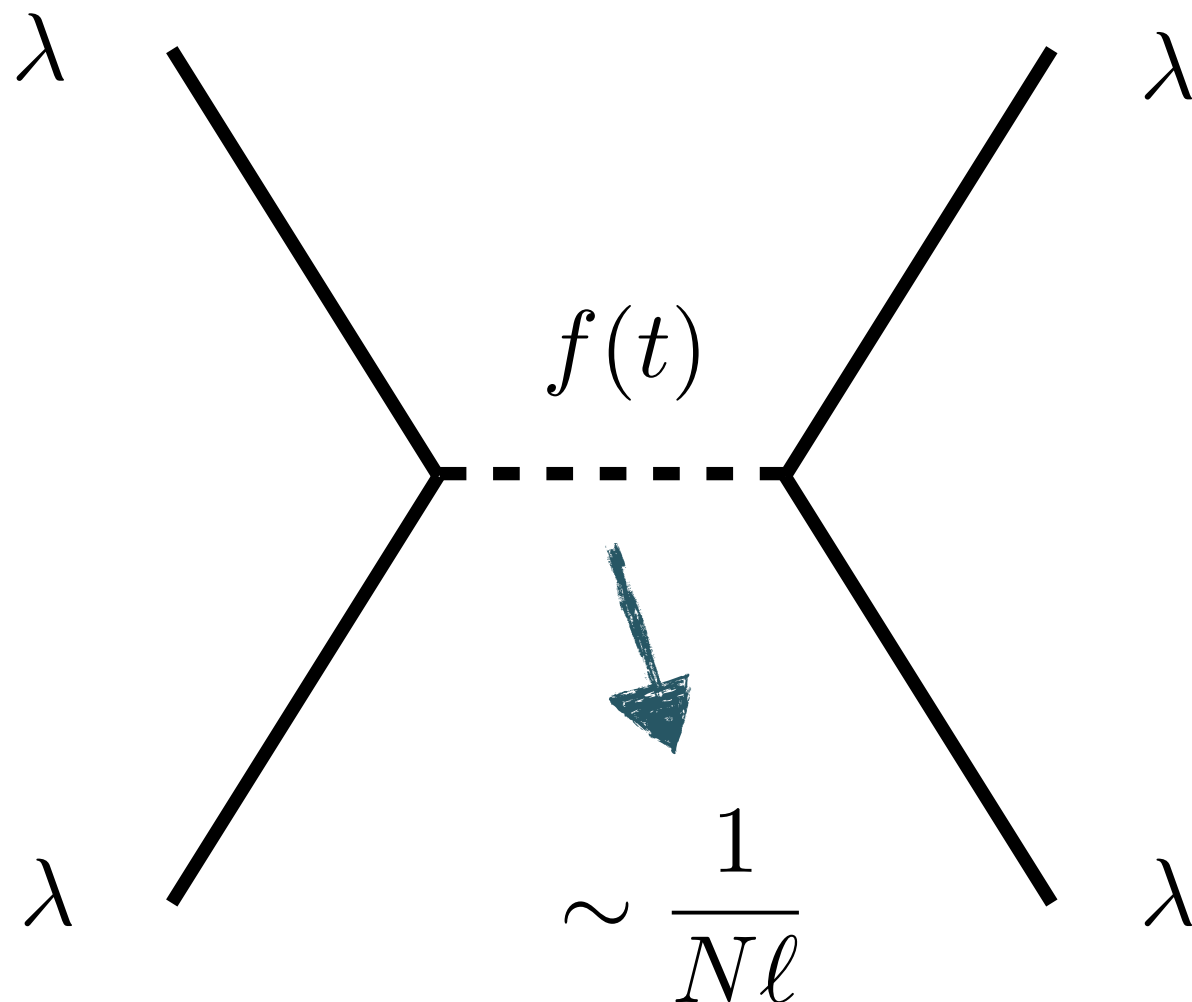
Lyapunov exponent in $\text{NAdS}_2/\text{NCFT}_1$

Compute Euclidean four-point function at tree-level:



Lyapunov exponent in NAdS₂/NCFT₁

Compute Euclidean four-point function at tree-level:



Analytically
continue to get
out-of-time-ordered
four-point function

$$\lambda_L = \frac{2\pi}{\beta}$$

Outline

1. Motivation
2. $\text{NAdS}_2/\text{NCFT}_1$
3. Chaos
- 4. Parting words**

A sigma model for AdS_3 gravity

WIP with Jordan Cotler (Stanford)

It is possible to rewrite pure classical AdS_3 gravity as an unconventional sigma model from $\mathcal{M}_2 \rightarrow \mathcal{M}_2$

A sigma model for AdS₃ gravity

WIP with Jordan Cotler (Stanford)

The action may be regarded as an action for an “exact” hydrodynamics whose gradient expansion truncates at second order in derivatives

$$S = \frac{c}{24\pi} \int d^2\sigma \sqrt{-g} \left(4\pi^2 T^2 - \frac{(\partial T)^2}{T^2} - \ln T R \right)$$

A sigma model for AdS₃ gravity

WIP with Jordan Cotler (Stanford)

$$S = \frac{c}{24\pi} \int d^2\sigma \sqrt{-g} \left(4\pi^2 T^2 - \frac{(\partial T)^2}{T^2} - \ln T R \right)$$

Torus partition function?

Coupling to matter and Virasoro blocks?

Stay tuned!

Thank you!