

# Sharp Boundaries for the Swampland

w/ Caron-Huot, Mazur, Rastelli

Q: Which EFT's can be UV-completed?

Not everything goes:

- w/out gravity

$$\mathcal{L} \supseteq g_2 (\partial\phi)^4 + g_3 \partial^6 \phi^4 + \dots$$

$$\mathcal{M}_{\text{low}}(s,u) \supseteq g_2 (s^2 + t^2 + u^2) + g_3 stu + \dots$$

$$\text{UV complete} \Rightarrow g_2 \geq 0 \quad [\text{Adams et al. '06}]$$

2 arguments:

- dispersion relation
- causality in nontrivial background

Recently: [Arkani-Hamed, Huang, Huang '20]

[Bellazzini, Miao, Rattazzi, Rombau, Riva '20]

[Tolley, Wang, Zhou '20]

[Caron-Huot, Van-Duong '20]

...

$$\text{sharp bounds: } -\frac{\#}{M^2} \leq \frac{g_3}{g_2} \leq \frac{\#}{M^2}$$

↑ UV cutoff

Slogan: dimensional analysis is a theorem

- with gravity

[Caroniu, Edelstein, Maldacena, Zhiboedov '14], ...

$$A_{ggg} = \sqrt{32\pi G} (A_R + \alpha_2 A_{R^2} + \alpha_4 A_{R^3})$$

$$\alpha_2 \sim \frac{1}{M^2}, \quad \alpha_4 \sim \frac{1}{M^4}$$

also: nonzero  $\alpha_2, \alpha_4 \Rightarrow \infty$  tower of massive spinning particles

- using causality
- only parametric bounds

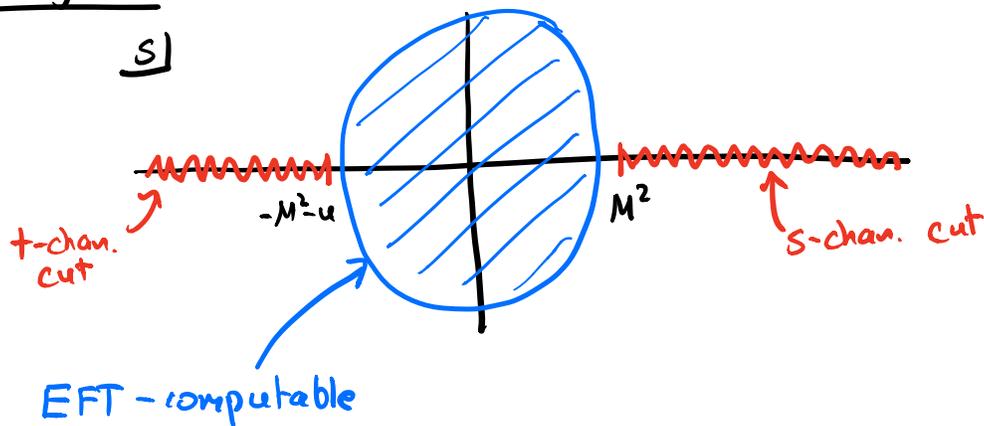
This talk: Sharp bounds with gravity

Key Assumptions + Dispersion Review

Consider  $\phi\phi \rightarrow \phi\phi$   $\mathcal{M}(s, u)$

- 1. Analyticity
- 2. Unitarity
- 3. Regge bound

1. Analyticity



simplicity: tree-level

$$\mathcal{M}_{\text{low}}(s, u) = 8\pi G \left( \frac{st}{u} + \frac{su}{t} + \frac{tu}{s} \right) - \lambda_3^2 \left( \frac{1}{s} + \frac{1}{t} + \frac{1}{u} \right) - \lambda_4$$

$$+ g_2 (s^2 + t^2 + u^2) + g_3 stu + g_4 (s^2 + t^2 + u^2)^2 + \dots$$

- could include loops
  - nonlinear map between  $\mathcal{L}$  and  $\mathcal{M}_{low}$
  - we bound obs. linear in  $\mathcal{M}_{low}$

## 2. Unitarity

partial wave decomp:

$$M(s, u) = s^{\frac{4-D}{2}} \sum_{J \text{ even}} a_J c_J(s) P_J\left(1 + \frac{2u}{s}\right)$$

$\uparrow$  Gegenbauer (Legendre in  $D=4$ )       $\uparrow$   $\cos \theta$  scattering angle

$$S S^\dagger = 1, \quad S = 1 + iM$$

$$\Rightarrow |1 + i c_J(s)|^2 \leq 1$$

$$p_J(s) \equiv \text{Im } c_J(s)$$

$$\Rightarrow \underbrace{0 \leq p_J(s) \leq 2}_{\text{this talk}} \quad (s > 0, J \text{ even})$$

## 3. Bounded Regge Growth (spin 2)

$$\lim_{|s| \rightarrow \infty} \frac{M(s, u)}{s^2} = 0 \quad (u < 0)$$

define spin- $k$ -subtracted disp. relation

$$C_{k,u}: 0 = \int_0^\infty \frac{ds}{s} \frac{M(s,u)}{(s(s+u))^{k/2}}$$

boundedness assumption  $\Rightarrow C_{k,u}$  converges for  $k=2, 4, \dots$

### Comments

- Satisfied in weakly-coupled string theory

$$M \sim s^{2+\alpha' u}$$

$\uparrow$  Reggeization of graviton

-  $M(s,u) \leq O(s^2)$  is the

"Classical Regge Growth" conjecture

[Caroniu, Edelstein, Maldacena, Zhiboedov '14]

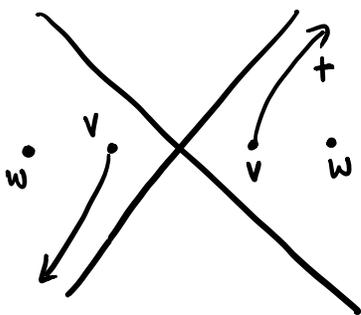
[Maldacena, Shenker, Stanford '15]

[Chandarkar, Chowdhury, Kundu, Minwalla '21]

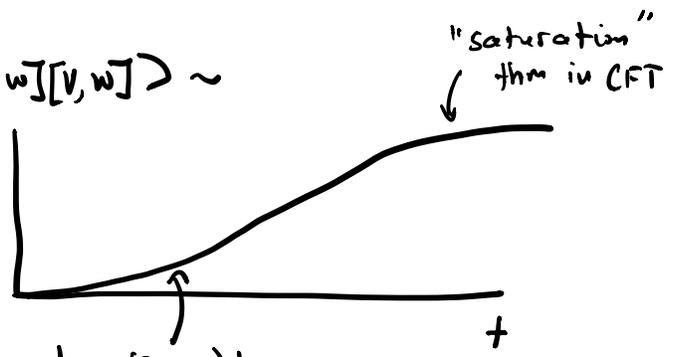
We actually need something stronger

$\Rightarrow$  Rigorously provable in AdS

Aside: AdS/CFT + chaos



$$\langle [v, w][v, w] \rangle \sim$$



$$\frac{1}{N^2} e^{(J_{\text{eff}}-1)t}$$

chaos bound:  $J_{\text{eff}} \leq 2$

Spin- $k$  dispersion relation

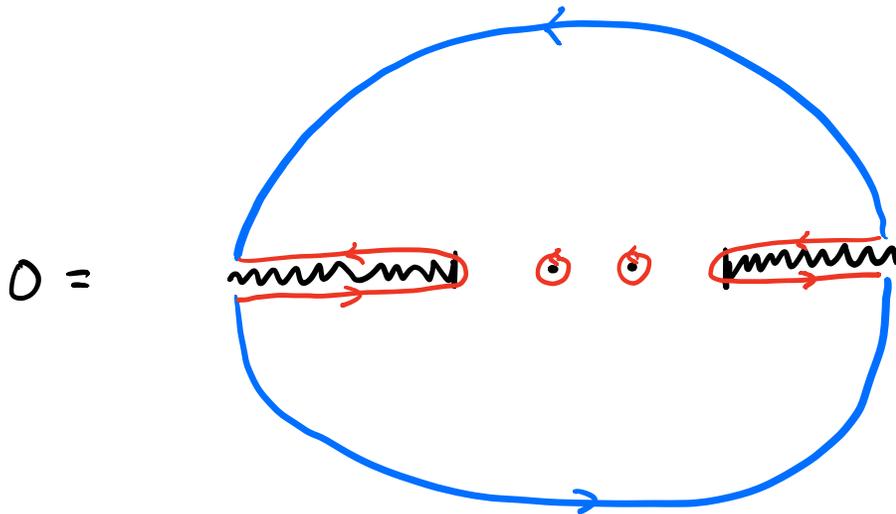
$$\sim (\dots) = \int_0^\infty dt e^{-(k-1)t} \langle [V, W][V, W] \rangle$$

2 examples of convergence:

- $J_{\text{eff}} < 2 \Rightarrow \int_0^\infty dt e^{(J_{\text{eff}}-k)t} < \infty$   
(weakly-coupled string theory)
- saturation  $\Rightarrow \int_0^\infty dt e^{(1-k)t} < \infty$

(Upcoming work [w/ Caron-Huot, Mazur, Rastelli])  
flat-space limit of CFT dispersion relations  
 $\Rightarrow$  this talk

## Dispersion Relations: Applications



$$\mathcal{L}_{2,u} = \frac{8\pi G}{-u} + 2g_2 - g_3 u + 8g_4 u^2 + \dots = \left\langle \frac{(2m^2 + u) P_J \left(1 + \frac{2u}{m^2}\right)}{m^2 (m^2 + u)^2} \right\rangle$$

$$\mathcal{C}_{4,u} : 4g_4 + \dots = \left\langle \frac{(2m^2+u) P_J(1+\frac{2u}{m^2})}{m^4 (m^2+u)^3} \right\rangle$$

$$\text{where } \langle \dots \rangle = \frac{1}{\pi} \sum_J n_J \int_{M^2}^{\infty} \frac{dm^2}{m^2} m^{4-D} \rho_J(m^2) (\dots)$$

Review: Adams et al.  $8\pi G = 0$

$$\mathcal{C}_{2,u=0} : 2g_2 = \left\langle \frac{2}{m^4} \right\rangle \geq 0$$

↑ forward limit  $\cos\theta = 1$   $P_J(1) = 1$

Now turn on G

$$\text{problem: } \left. \frac{8\pi G}{-u} \right|_{u=0} = \infty$$

useful to view problem in impact param. space

$$u = -p^2$$

$$p = |\vec{p}|$$

$$b = |\vec{b}|$$

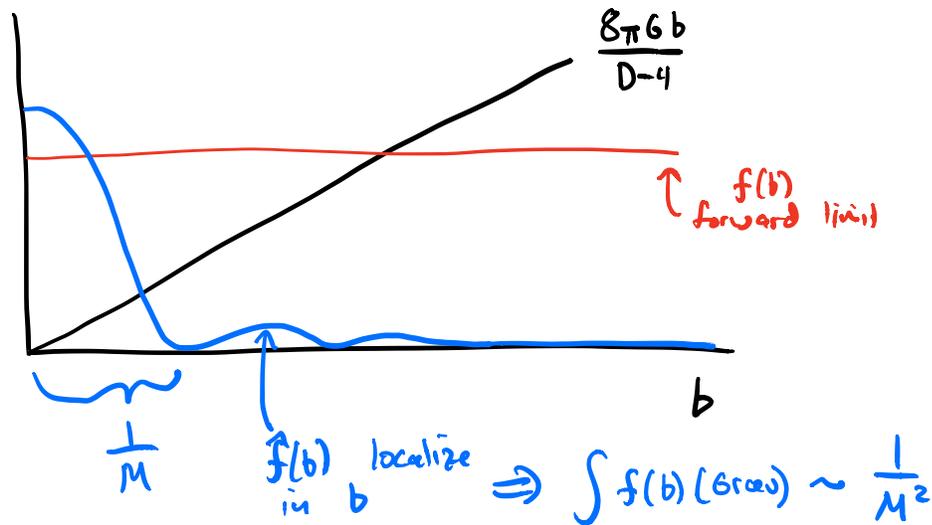
↑  
transv. momentum transfer

↑  
impact param

$$\int_0^{\infty} dp f(p) \frac{8\pi G}{p^2} = \int_0^{\infty} db \hat{f}(b) \frac{8\pi G b}{D-4}$$

$$\hat{f}(b) \equiv \int d^{D-2} \vec{p} e^{i\vec{b} \cdot \vec{p}} \frac{f(p)}{p^{D-3} \text{ vol } S^{D-3}}$$

forward limit :  $f(p) = \delta(p)$ ,  $\hat{f}(b) = 1$



localize in  $b$ : need  $|u| \sim M^2$

Worry: EFT series breaks down

Solution: "improved" sum rules

$$\mathcal{C}_{2,u}^{\text{improved}} = C_{2,u} - \left( \underset{\substack{\uparrow \\ \text{cancels } 8g_4 u^2}}{2u^2 C_{4,0}} + u^3 \underset{\substack{\uparrow \\ \text{cancels } g_6}}{C'_{4,0}} + \dots \right)$$

$$\mathcal{C}_{2,u}^{\text{imp.}} : \frac{8\pi G}{-u} + 2g_2 - g_3 u = \left\langle \frac{(2m^2+u) P_J(1+\frac{2u}{m^2})}{m^2(m^2+u)^2} - \frac{u^2}{mb} \left( \frac{(4m^2+3u) P_J(1)}{(m^2+u)^2} + \frac{P_J'(1)}{mu} \right) \right\rangle$$

$$= \langle C_{2,u}^{\text{imp.}} [m^2, J] \rangle$$

valid for  $u = -p^2$ ,  $p \in (0, M)$

Now: linear programming bounds

search for  $f(p)$  s.t.

$$\int_0^M dp f(p) \mathcal{L}_{2,-p^2}^{\text{imp}} [m^2, \mathcal{J}] \geq 0 \quad \text{for all } m > M, \mathcal{J} = 0, 2, 4, \dots$$

$$\Rightarrow \int_0^M dp f(p) \left[ \frac{8\pi G}{p^2} + 2g_2 + g_3 p^2 \right] \geq 0$$

(can choose  $f(p)$  to optimize different things)

Intuition

Scaling limit  $m \rightarrow \infty$  w/

$$b_{\text{eff}} = \frac{2\mathcal{J}}{m} \text{ held fixed}$$

$$\int_0^M dp f(p) \mathcal{L}_{2,-p^2}^{\text{imp}} [m^2, \mathcal{J}] \sim \frac{\#}{m^4} \hat{f}(b_{\text{eff}})$$

$\Rightarrow$  bounds from - positive  $\hat{f}(b)$   
- c-ct. support  $f(p)$

$$\left[ \text{example: } \int_0^M \frac{dp}{M} \left(1 - \frac{p}{M}\right) \cos(pb) = \frac{1 - \cos(Mb)}{(Mb)^2} \geq 0 \right]$$

In practice

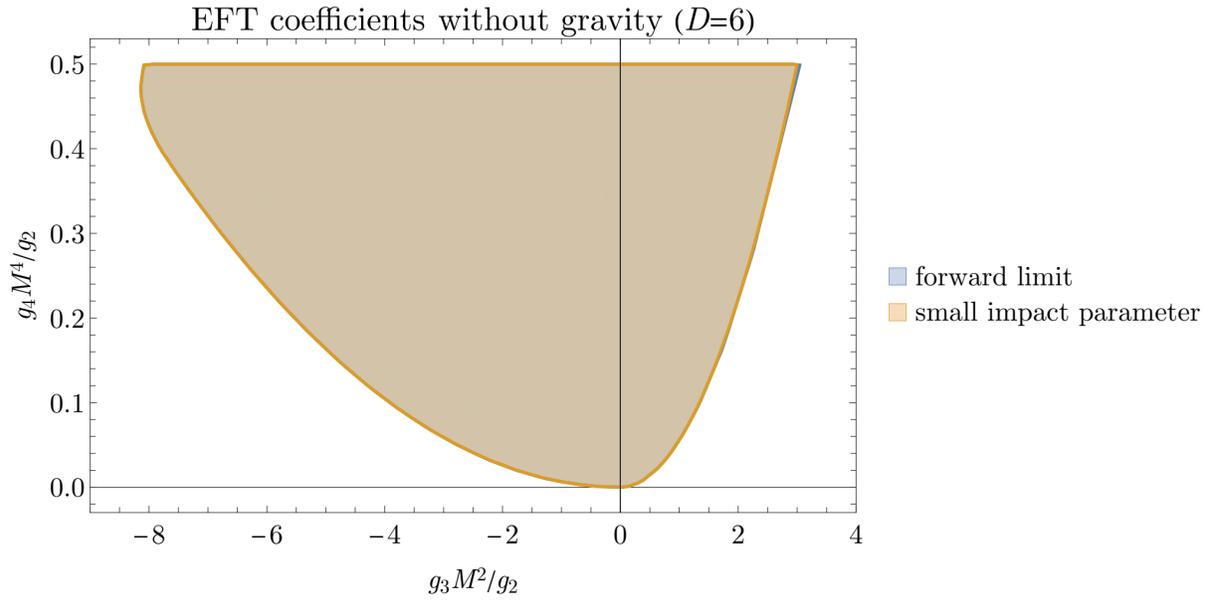
$$\text{write: } f(p) = \sum_n a_n p^n$$

impose positivity: - for  $\mathcal{J} = 0, 2, \dots, \mathcal{J}_{\text{max}}$   
- also in  $b$ -space

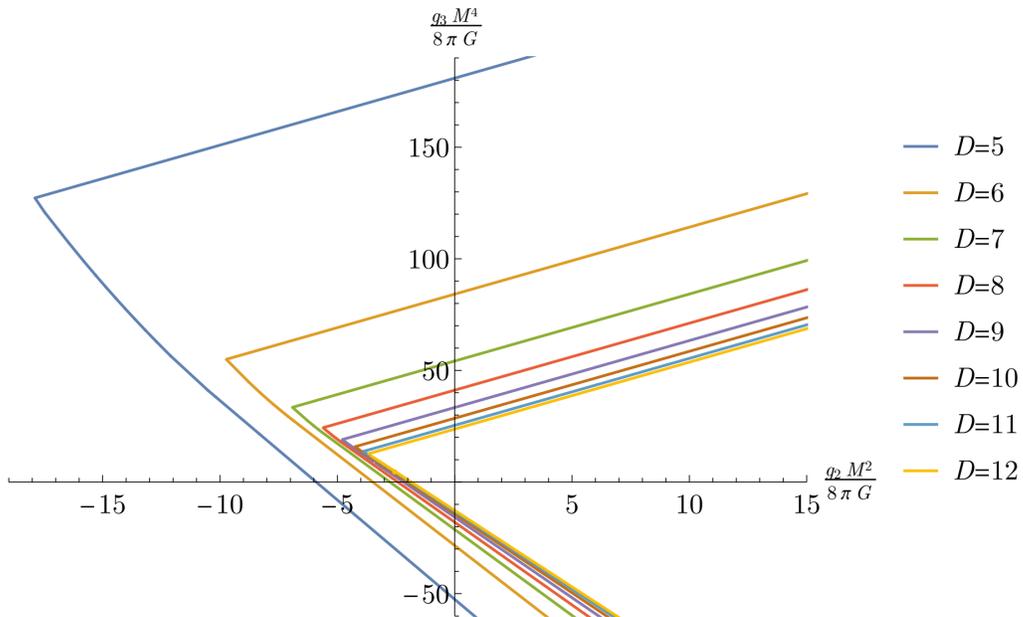
solve w/ SDPB [DSD '15], [Landry, DSD '19]

# Results

no gravity: reproduces same results as forward limit



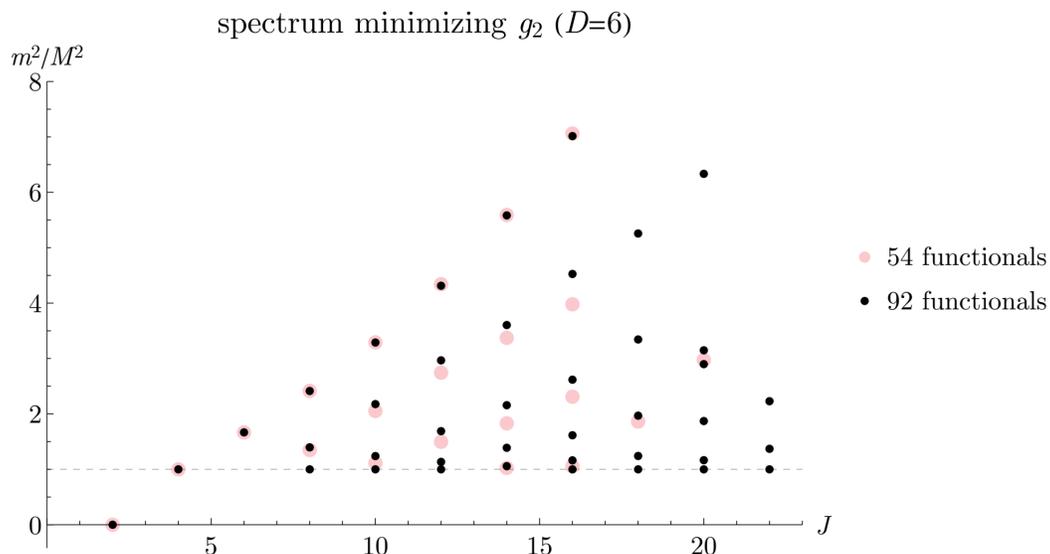
New: bounds with gravity



Note:  $g_2$  can be negative, but has a lower bound.

(intuition: negative  $g_2$  gives time advance, but this is ok because gravity gives a time delay)

"Extremal" spectrum



(does not look like string theory)

D=4 + IR divergences

$$\int db \hat{\mathcal{F}}(b) \frac{8\pi G b}{D-4}$$

↑ impact param amplitude for grav. diverges when  $D=4$

$$g_2 \approx -\frac{8\pi G}{M^2} \times 17 \log(1.7 M b_{\max}) \quad (D=4)$$

↑  
IR cutoff