Probing Lepton Flavor Violation at Future Electron Positron Colliders

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based on arXiv:2305.03869 with Pankaj Munbodh and Talise Oh

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Future Colliders



Karl Jacobs (ECFA chair) @ 2nd ECFA meeting on e^+e^- Higgs, electroweak, and top factories Oct 11-13, 2023, Paestum, Italy

 In the SM, charged lepton flavor violation is suppressed by the tiny neutrino mass splittings

e.g.
$$\mathsf{BR}(\mu \to 3e) \sim \mathsf{BR}(\mu \to e \nu_e \nu_\mu) \left| \frac{g^2}{16\pi^2} \frac{\Delta m_\nu^2}{m_W^2} \right|^2 \sim 10^{-50}$$

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- ► Can search for lepton flavor violation in many different ways:
- 1) At low energies in lepton or hadron decays: $\mu \rightarrow e\gamma$, $B_s \rightarrow \tau\mu$, ...

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- 3) At high energies in non-resonant production: $e^+e^- \rightarrow \tau \mu$, ...

New Physics Sensitivity of LFV at Low Energies

► Generic scaling of a new physics effect with the flavor changing coupling g_{NP} and the new physics scale Λ_{NP}

$$rac{{\sf BR}(\mu o 3e)}{{\sf BR}(\mu o e
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▶ For O(1) couplings, this corresponds to new physics scales of

 $\Lambda_{NP} \gtrsim 100 \text{ TeV}$ for muons $\Lambda_{NP} \gtrsim 10 \text{ TeV}$ for taus

New Physics Sensitivity of Heavy Resonance Decays

 Consider LFV decays of the Z boson, the Higgs, the top in the presence of generic new physics

$$\begin{split} \frac{\mathsf{BR}(Z \to \mu e)}{\mathsf{BR}(Z \to \mu \mu)} &\sim g_{\mathsf{NP}}^2 \left(\frac{v}{\Lambda_{\mathsf{NP}}}\right)^4 \;, \quad \frac{\mathsf{BR}(H \to \tau \mu)}{\mathsf{BR}(H \to \tau \tau)} \sim g_{\mathsf{NP}}^2 \left(\frac{v}{\Lambda_{\mathsf{NP}}}\right)^4 \\ & \frac{\mathsf{BR}(t \to c \mu e)}{\mathsf{BR}(t \to W b)} \sim \frac{g_{\mathsf{NP}}^2}{\mathsf{16}\pi^2} \left(\frac{v}{\Lambda_{\mathsf{NP}}}\right)^4 \end{split}$$

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$$\frac{\mathsf{BR}(Z \to \mu e)}{\mathsf{BR}(Z \to \mu \mu)} \sim g_{\mathsf{NP}}^2 \left(\frac{v}{\Lambda_{\mathsf{NP}}}\right)^4 , \quad \frac{\mathsf{BR}(H \to \tau \mu)}{\mathsf{BR}(H \to \tau \tau)} \sim g_{\mathsf{NP}}^2 \left(\frac{v}{\Lambda_{\mathsf{NP}}}\right)^4 \\ \frac{\mathsf{BR}(t \to c \mu e)}{\mathsf{BR}(t \to W b)} \sim \frac{g_{\mathsf{NP}}^2}{16\pi^2} \left(\frac{v}{\Lambda_{\mathsf{NP}}}\right)^4$$

- ► Same dependence on new physics as the low energy probes, but typically much less *Z*, Higgs, top available in experiments.
- Note: these are extremely generic/naive expectations; situation can be very different in concrete models.

[for a review see WA, Caillol, Dam, Xella, Zhang 2205.10576]

▶ Results from the LHC: ATLAS (139 fb⁻¹)

Phys.Rev.Lett. 127 (2022) 271801; Nature Phys. 17 (2021) 7, 819-825; ATLAS-CONF-2021-042

 ${\sf BR}(Z o \mu e) < 3.04 imes 10^{-7} \ {\sf BR}(Z o au e) < 5.0 imes 10^{-6} \ {\sf BR}(Z o au \mu) < 6.5 imes 10^{-6}$

- ► Slightly better than LEP bounds for all decay modes.
- In all searches there are backgrounds ⇒ expect sensitivities to improve with √L, i.e. ~ factor of 5 at the HL-LHC.

Expected Sensitivities at Proposed Z Pole Machines

based on FCC-ee study Dam 1811.09408 (see also the FCC-ee whitepaper 2203.06520)

- background from Z → ττ → μνν eνν is under control. Momentum resolution of 10⁻³ and Z mass constraint implies background rate of ~ 10⁻¹¹.
- ▶ main background: $Z \rightarrow \mu\mu$ where one muon suffers from "catastrophic" bremsstrahlung and is identified as electron.
- ► mis-id probability $\sim 10^{-7}$ limits the sensitivity to BR($Z \rightarrow \mu e$) $\sim 10^{-8}$.
- With improved e/µ separation (dE/dx) might be able to go down to BR(Z → µe) ~ 10⁻¹⁰.

 $Z \rightarrow \mu e$

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- With improved *e*/µ separation (*dE*/*dx*) might be able to go down to BR(Z → µe) ~ 10⁻¹⁰.

$$\begin{array}{ccc} Z \rightarrow \tau e & & \blacktriangleright & \mathsf{mi} \\ \text{and} & & \flat & \mathsf{ba} \\ Z \rightarrow \tau \mu & & & \vdots \end{array}$$

 $Z \rightarrow \mu e$

• minimize τ vs μ , e mis-id \rightarrow focus on hadronic taus

• background from
$$Z \rightarrow \tau_{had} \tau \rightarrow \tau_{had} \ell \nu \nu$$

▶ limits sensitivity to $BR(Z \rightarrow \tau \ell) \sim 10^{-9}$

LFV Z Decays in the EFT Framework

 Parameterize New Physics in a systematic and controlled way: in terms of dim-6 operators of the SMEFT

dipoles

$$\mathcal{O}_{dW} = (\bar{\ell}\sigma^{\mu\nu}\tau^a P_R\ell')H \ W^a_{\mu\nu}$$

$$\mathcal{O}_{dB} = (\bar{\ell} \sigma^{\mu\nu} P_R \ell') H \ B_{\mu\nu}$$



Higgs currents

$$\mathcal{O}_{hl}^{(3)} = (H^{\dagger}i\overleftrightarrow{\mathsf{D}}_{\mu}^{a}H)(\bar{\ell}\gamma^{\mu}\tau^{a}P_{L}\ell')$$
$$\tilde{\mathcal{O}}_{hl}^{(1)} = (H^{\dagger}i\overleftrightarrow{\mathsf{D}}_{\mu}H)(\bar{\ell}\gamma^{\mu}P_{L}\ell')$$
$$\mathcal{O}_{he} = (H^{\dagger}i\overleftrightarrow{\mathsf{D}}_{\mu}H)(\bar{\ell}\gamma^{\mu}P_{R}\ell')$$



Comparison with Low Energy Probes

- ► Many flavor violating low energy processes will be affected as well.
- Severe indirect constraints on $Z \rightarrow \mu e$ from $\mu \rightarrow e\gamma$, $\mu \rightarrow 3e$, $\mu \rightarrow e$ conversion (barring accidental cancellations).



Calibbi, Marcano, Roy 2107.10273

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Complementary sensitivity in the case of taus.



$$\frac{\sigma(\mathbf{e}^+\mathbf{e}^- \to \tau\mu)}{\sigma(\mathbf{e}^+\mathbf{e}^- \to \tau^+\tau^-)} \sim$$

$$\frac{\sigma(\boldsymbol{e}^+\boldsymbol{e}^- \to \tau\mu)}{\sigma(\boldsymbol{e}^+\boldsymbol{e}^- \to \tau^+\tau^-)} \sim g_{\rm NP}^2\left(\frac{\boldsymbol{v}^4}{\Lambda_{\rm NP}^4}\right),$$

$$\frac{\sigma(\boldsymbol{e}^{+}\boldsymbol{e}^{-} \to \tau \mu)}{\sigma(\boldsymbol{e}^{+}\boldsymbol{e}^{-} \to \tau^{+}\tau^{-})} \sim g_{\rm NP}^{2} \left(\frac{\boldsymbol{v}^{4}}{\Lambda_{\rm NP}^{4}}\right), \ g_{\rm NP}^{2} \left(\frac{\boldsymbol{s}\boldsymbol{v}^{2}}{\Lambda_{\rm NP}^{4}}\right),$$

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- For some operators one will have enhanced sensitivity at high energies. (Assuming one does not resolve the higher dimensional operators.)
- ▶ How sensitive is one to $\tau\mu$ production at future e^+e^- colliders?

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- For some operators one will have enhanced sensitivity at high energies. (Assuming one does not resolve the higher dimensional operators.)
- ► How sensitive is one to $\tau\mu$ production at future e^+e^- colliders?
- In WA, Munbodh, Oh 2305.03869 we show that high-energy runs of FCC-ee/CEPC have sensitivity that is comparable and complementary to other probes.

Systematic SMEFT Parameterization of New Physics

dipoles

$$\mathcal{O}_{dW} = (\bar{\tau}\sigma^{\alpha\beta}T^{a}P_{R}\mu)H \ W^{a}_{\alpha\beta}$$
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Higgs currents

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4-fermion contact interactions

$$\mathcal{O}_{\ell\ell} = (\bar{e}\gamma^{\alpha}P_{L}e)(\bar{\tau}\gamma_{\alpha}P_{L}\mu)$$
$$\mathcal{O}_{ee} = (\bar{e}\gamma^{\alpha}P_{R}e)(\bar{\tau}\gamma_{\alpha}P_{R}\mu)$$
$$\mathcal{O}_{\ell e} = (\bar{e}\gamma^{\alpha}P_{L}e)(\bar{\tau}\gamma_{\alpha}P_{R}\mu)$$
$$\mathcal{O}_{e\ell} = (\bar{e}\gamma^{\alpha}P_{R}e)(\bar{\tau}\gamma_{\alpha}P_{L}\mu)$$







Dependence on the Center of Mass Energy



WA, Munbodh, Oh 2305.03869 (in the plot $\Lambda_{NP} = 3$ TeV. $C_i = 1$)

- τµ production increases linearly with s for 4-fermion operators
- *τ* μ production is flat in
 s for dipole operators
- τμ production falls like 1/s for Higgs current operators
- resonance at $s = m_Z^2$ if *Z*-mediated

Signal and Most Important Background

signal: $e^+e^- \rightarrow \tau \mu$

bkg: $e^+e^- \rightarrow \tau^+\tau^- \rightarrow \tau\mu\nu\nu$

- Signal is a sharp peak at $x = p_{\mu}/p_{\text{beam}} = 1$
- Background is a smooth distribution with $x \leq 1$
- Width of the signal peak and spread of background to x > 1 is determined by the beam energy spread and the muon momentum resolution.



(study on the Z peak)

Impact of initial state radiation? (work in progress with Munbodh)

Another Background at High Energies?

$e^+e^- ightarrow W^+W^- ightarrow au\mu u u$

- Muon momentum does not extend all the way to x = 1
- Decay kinematics is such that

$$x < \frac{1}{2} \left(1 + \sqrt{1 - \frac{4m_W^2}{s}} \right) < 1$$

• e.g. for $\sqrt{s} = 240$ GeV one has $x \lesssim 0.87$

\Rightarrow this background is not an issue.

Existing Constraints from LEP



WA, Munbodh, Oh 2305.03869

- ► LEP has searched for $e^+e^- \rightarrow \tau\mu$ at the Z pole (e.g. OPAL Z.Phys.C 67 (1995) 555-564) and at $\sqrt{s} \sim 200 \text{ GeV}$ (OPAL PLB 519, (2001) 23-32).
- Z pole search mainly sensitive to the Higgs current operators.
- ► High √s search mainly sensitive to 4-fermion operators.
- ► LEP searches have sensitivity comparable to $Z \rightarrow \tau \mu$ at the LHC, but cannot compete with tau decays.

Projections for FCC-ee

machine and detector parameters from FCC-ee CDR vol. 2, 1909.12245, 2107.02686, 2203.06520

$\sqrt{s} [\text{GeV}]$	$\mathcal{L}_{int} \ [ab^{-1}]$	$\frac{\delta\sqrt{s}}{\sqrt{s}} \ [10^{-3}]$	$\frac{\delta p_T}{p_T} \left[10^{-3} \right]$	$\epsilon_{\rm bkg}^{x_c} \ [10^{-6}]$	$N_{\rm bkg}$	$\sigma~[{\rm ab}]$
91.2 (Z-pole)	75	0.93	1.35	1.55	9700 ± 100	45
$87.7 \ (\text{off-peak})$	37.5	0.93	1.33	1.46	520 ± 20	21
93.9 (off-peak)	37.5	0.93	1.37	1.59	930 ± 30	28
125 (H)	20	0.03	1.60	1.44	12 ± 3	8
$160 \; (WW)$	12	0.93	1.89	2.44	6 ± 2	10
240~(ZH)	5	1.17	2.60	4.39	2 ± 1	18
$365~(t\bar{t})$	1.5	1.32	3.78	8.61	0.5 ± 0.7	50

- Estimate background efficiency by imposing a cut x > 1. (could be further optimized)
- Expect sizable background on the Z-peak, very few background events at higher energies.
- ▶ Can achieve sensitivity to $e^+e^- \rightarrow \tau \mu$ cross sections of O(10 ab).

Projections for CEPC

machine and detector parameters from 1809.00285, 1811.10545, 2203.09451, 2205.08553

$\sqrt{s} [\text{GeV}]$	$\mathcal{L}_{int} \ [ab^{-1}]$	$\frac{\delta\sqrt{s}}{\sqrt{s}}~[10^{-3}]$	$\frac{\delta p_T}{p_T} \ [10^{-3}]$	$\epsilon_{\rm bkg}^{x_c} \ [10^{-6}]$	$N_{\rm bkg}$	$\sigma~[{\rm ab}]$
91.2 (Z-pole)	50	0.92	1.35	1.53	6400 ± 80	55
87.7 (off-peak)	25	0.92	1.33	1.46	350 ± 20	27
93.9 (off-peak)	25	0.92	1.37	1.59	620 ± 25	35
160~(WW)	6	0.99	1.89	2.49	3 ± 2	17
240~(ZH)	20	1.20	2.60	4.42	7 ± 3	6.6
$360~(t\bar{t})$	1	1.41	3.74	8.61	0.3 ± 0.5	72

- Estimate background efficiency by imposing a cut x > 1. (could be further optimized)
- Expect sizable background on the Z-peak, very few background events at higher energies.
- ► Can achieve sensitivity to $e^+e^- \rightarrow \tau \mu$ cross sections of $\mathcal{O}(10 \text{ ab})$.

Complementarity of Different Observables (FCC-ee)



WA, Munbodh, Oh 2305.03869

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Summary of Generic Sensitivities



WA, Munbodh, Oh 2305.03869

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If a signal is seen at one √s:
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- If a signal is seen at one √s:
 ⇒ look at different √s to identify the operator class (dipole, Higgs current, 4-fermion)
- The signal can be further characterized by angular distributions (θ = angle between the beam axis and the outgoing muon) and CP asymmetries (τ⁺μ[−] vs. τ[−]μ⁺)

$$\frac{1}{\sigma_{\text{tot}}} \frac{d(\sigma + \bar{\sigma})}{d\cos\theta} = \frac{3}{8} (1 - F_D) (1 + \cos^2\theta) + A_{\text{FB}} \cos\theta + \frac{3}{4} F_D \sin^2\theta ,$$
$$\frac{1}{\sigma_{\text{tot}}} \frac{d(\sigma - \bar{\sigma})}{d\cos\theta} = \frac{3}{8} (A^{\text{CP}} - F_D^{\text{CP}}) (1 + \cos^2\theta) + A_{\text{FB}}^{\text{CP}} \cos\theta + \frac{3}{4} F_D^{\text{CP}} \sin^2\theta ,$$

► For a sufficiently large signal, it might be possible to significantly narrow down the chirality structure of the operator that is responsible for $e^+e^- \rightarrow \tau \mu$

- ► Non-resonant $e^+e^- \rightarrow \tau \mu$ offers interesting opportunities to probe lepton flavor violation at FCC-ee/CEPC.
- Different LFV operators show characteristic dependence on the center of mass energy.
- Estimated sensitivity rivals the one from rare tau decays.
- Most relevant machine/detector parameters: beam energy spread and muon momentum resolution.
- Linear colliders are also interesting: higher center of mass energy and polarized beams.