HOLOGRAPHIC ENTANGLEMENT

Veronika Hubeny



Physics Department & center for Quantum Mathematics and Physics



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based on earlier works w/ {M. Headrick, A. Lawrence, H. Maxfield, M. Rangamani, T. Takayanagi, E. Tonni} & on work in progress w/ M. Headrick

OUTLINE

- Context
 - Background: AdS/CFT
 - Motivation: covariantize to elucidate the correspondence
- Holographic Entanglement
 - Entanglement entropy
 - RT & HRT
- Recasting holographic entanglement
 - Bit threads
 - Covariant bit threads
- Summary

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Gauge/gravity duality



• Since the two descriptions live in different number of dimensions, we call such a correspondence holographic.

Gauge/gravity duality

* better analogy: stereogram...



...but infinitely more complicated

Radial direction

What boundary quantity encodes the extra bulk direction?

• hints from spatial geometry of AdS:



Radial direction

What boundary quantity encodes the extra bulk direction?



• hints from spatial geometry of AdS:

The radial bulk direction comes from a scale size on boundary:

- points near the boundary \leftrightarrow small arcs
- points further in the bulk \leftrightarrow larger arcs
- same point represented by multiple arcs

- Provides useful intuition:
 e.g. bulk particle falling due to gravity
 ... falling into black hole
 - ↔ boundary excitation spreading outward↔ ... thermalizing

Applications of gauge/gravity duality

String theory (gravity) ↔ "in bulk" = higher dimensions

describes gravitating systems, e.g. black holes



field theory (no gravity) "on boundary" = lower dimensions describes experimentally accessible systems



Invaluable tool to:

- Study strongly interacting field theory (hard, but describes many systems) by working with higher-dimensional gravity on AdS (easy).
- Study quantum gravity in AdS (hard, but needed to understand spacetime) by using the field theory (easy for certain things)

Pre-requisite:

We need to understand the AdS/CFT dictionary...

- How does bulk spacetime emerge from the CFT?
 - Which CFT quantities give the bulk metric?
 - What determines bulk dynamics (Einstein's eq.)?
 - How does one recover a local bulk operator from CFT quantities?
- What part of bulk can we recover from a restricted CFT info?
 - What bulk region does a CFT state (at a given instant in time) encode?
 - What bulk region does a spatial subregion of CFT state encode?
- (How) does the CFT "see" inside a black hole?
 - Does it unitarily describe black hole formation & evaporation process?
 - How does it resolve curvature singularities?

Main message: using GR technology goes a long way...

Motivation

- Elucidate holography
 - Fundamental nature of spacetime & its relation to entanglement
 - Structure/characterization of CFTs (& states) w/ gravity dual
- Start w/ situations with large amount of symmetry (e.g. pure AdS)
 - Explicit calculations possible, can obtain analytical expressions
 - Use these to guess duality relations → entry in gauge/gravity dictionary
- But this has limitations
 - How to generalize? (e.g. time dependence)
 - Often symmetry brings degeneracy between logically distinct concepts
- Need to "covariantize"
 - Define a quantity which is purely geometrical (e.g. independent of any choice of coordinate systems) and fully general

Utility of covariant constructs

- Gives a general prescription
 - Definition of a quantity is equally robust on both sides of duality
 - Once beyond analytically tractable cases, might as well go for full generality (within the class of systems we want to consider, i.e. N = ∞)
- Time dependence interesting in its own right
 - Novel phenomena in out-of-equilibrium systems
 - New insight into the structure of the theory
- Breaks degeneracy between distinct constructs
 - Allows us to identify the true dual → underlying nature of the map
- Natural covariant constructs motivate new relations
 - Even if a given construct is not the sought dual, it eventually finds its use

Example: Holographic Entanglement Entropy

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Paths to Holographic Entanglement



Entanglement

Most non-classical manifestation of quantum mechanics

- "Best possible knowledge of a whole does not include best possible knowledge of its parts — and this is what keeps coming back to haunt us" [Schrodinger '35]
- New quantum resource for tasks which cannot be performed using classical resources [Bennet '98]
- Plays a central role in wide-ranging fields
 - quantum information (e.g. cryptography, teleportation, ...)
 - quantum many body systems
 - quantum field theory
- Hints at profound connections to geometry...

Entanglement Entropy (EE)

Suppose we only have access to a subsystem A of the full system = A + B. The amount of entanglement is characterized by Entanglement Entropy S_A :

• reduced density matrix $\rho_A = \text{Tr}_B |\psi\rangle\langle\psi|$ (more generally, for a mixed total state, $\rho_A = \text{Tr}_B \rho$)

• EE = von Neumann entropy $S_A = -{
m Tr}\,
ho_A\, \log
ho_A$

Defined if we can divide a quantum system into a subsystem A and its complement B, such that the Hilbert space decomposes:

$$\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$$

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 m Tr}\,
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- e.g. in local QFT:

A and B can be spatial regions, separated by a smooth entangling surface



The good news & the bad news

- But EE is hard to deal with...
 - non-local quantity, intricate & sensitive to environment
 - difficult to measure
 - difficult to calculate
 - ... especially in strongly-coupled quantum systems
- AdS/CFT to the rescue?
 - Is there a natural bulk dual of EE?
 (= ''Holographic EE'')



Yes! - described geometrically...

Holographic EE in static situations

Proposal [RT=Ryu & Takayanagi, '06] for static configurations:

In the bulk, entanglement entropy $S_{\mathcal{A}}$ for a boundary region \mathcal{A} is captured by the area of a minimal co-dimension-2 bulk surface \mathfrak{m} at constant t anchored on entangling surface $\partial \mathcal{A}$ & homologous to \mathcal{A}



$$S_{\mathcal{A}} = \min_{\substack{\partial \mathfrak{m} = \partial \mathcal{A}}} \frac{\operatorname{Area}(\mathfrak{m})}{4 \, G_N}$$

Remarks:

- Large body of evidence, culminating in [Lewkowycz & Maldacena '13]
- cf. black hole entropy...
- Minimal surface "hangs" into the bulk due to large distances near bdy.
- Note that both LHS and RHS are in fact infinite

Manifest properties of EE

E

- For pure states $S_{\mathcal{A}} = S_{\mathcal{A}^c}$
- Positivity: $S_{\mathcal{A}} \geq 0$
- Subadditivity: $S_{\mathcal{A}_1} + S_{\mathcal{A}_2} \geq S_{\mathcal{A}_1 \cup \mathcal{A}_2}$



• Implies positivity of mutual information: $I(A_1, A_2) = S_{A_1} + S_{A_2} - S_{A_1 \cup A_2}$

Proof of Strong Subadditivity

- strong subadditivity:
 - $S_{\mathcal{A}_1} + S_{\mathcal{A}_2} \geq S_{\mathcal{A}_1 \cup \mathcal{A}_2} + S_{\mathcal{A}_1 \cap \mathcal{A}_2}$
- proof in static configurations [Headrick & Takayanagi]



 $S_{\mathcal{A}_1} + S_{\mathcal{A}_2} = \alpha + \beta$

Proof of Strong Subadditivity

• strong subadditivity:

$$S_{\mathcal{A}_1} + S_{\mathcal{A}_2} \geq S_{\mathcal{A}_1 \cup \mathcal{A}_2} + S_{\mathcal{A}_1 \cap \mathcal{A}_2}$$

• proof in static configurations [Headrick & Takayanagi]



• Similarly prove monogamy of mutual information [Hayden, Headrick, Maloney] valid in holography but not in general: $S_A + S_B + S_C + S_{ABC} \leq S_{AB} + S_{BC} + S_{AC}$

Bulk dynamics from EE?

- We can in principle decode the bulk geometry from $\{S_A\}$ for a suitable set of A 's.
- But can we extract bulk dynamics more directly?
 - Use the strong subadditivity property of EE:

$$\delta^{2}_{\mathcal{A}} S_{\mathcal{A}} \sim \int_{\mathfrak{E}_{\mathcal{A}}} E_{ab} n^{a} n^{b} \geq 0$$
cf. Null Energy Condition

specific 2nd order variation of region

 proved at linearized level in 3-d, but conjectured to hold more generally...



[Bhattacharya,VH, Rangamani,Takayanagi, '14] cf. [Lashkari, Rabideau, Sabella-Garnier, Van Raamsdonk]

Spacetime from entanglement?

How does bulk spacetime emerge in the first place?

 Some connected spacetimes emerge as superpositions of disconnected spacetimes
 [Van Raamsdonk; Swingle]
 eg. eternal AdS black hole as thermofield double:



• Entanglement builds bridges: 'ER = EPR'





Einstein-Rosen bridge

Einstein-Podolsky-Rosen entanglement

Covariant Holographic EE

But the RT prescription is not well-defined outside the context of static configurations:

- In Lorentzian geometry, we can decrease the area arbitrarily by timelike deformations
- In time-dependent context, no natural notion of "const.t" slice...



In *time-dependent* situations, RT prescription must be covariantized:

Simplest candidate: [HRT = VH, Rangamani, Takayanagi '07]

minimal surface \mathfrak{m} at constant time

<u>extremal</u> surface \mathfrak{E} in the full bulk

This gives a well-defined quantity in any (arbitrarily time-dependent asymptotically AdS) spacetime \Rightarrow equally robust as in CFT



"Pf" in [Dong, Lewkowycz, Rangamani '16]

Curious features of EE:

- Extremal surfaces can have intricate behavior:
 - \mathfrak{E} can have discontinuous jumps under smooth variations of \mathcal{A}
 - → phase transitions in EE
 - ${\mathfrak E}$ can be topologically nontrivial even for simply-connected regions ${\mathcal A}$
- Holographic EE seems too local:
 - sharply-specified both on boundary and in bulk
 - but: \rightarrow we can reconstruct the bulk metric (modulo caveats) solely from the set $\{S_A\}$ for a suitable set of $\{A\}$
- Holographic EE seems too non-local:
 - global minimization condition + homology constraint makes S_A sensitive to arbitrarily distant regions in the bulk...

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Covariant re-formulations

- Covariance is pre-requisite to construct being physically meaningful, but it need not be unique
 - Distinct geometrical formulations can turn out equivalent
- This redundancy is useful
 - Each formulation can have its own advantages
 - e.g. different properties may be manifest in different formulations (cf. gauge / coordinate choice)
 - Re-formulation can reveal deeper relations (cf. ER=EPR [Maldacena, Susskind])

Covariant re-formulations of HEE

- $\mathfrak{E} = Extremal surface$
 - (relatively) easy to find
 - minimal set of ingredients required in specification
 - need to include homology constraint as extra requirement
- Φ = Surface with zero null expansions
 - cf. light sheet construction & covariant entropy bound [Bousso, '99]: Bulk entropy through light sheet of surface $\sigma \leq Area(\sigma)/4$
 - $\Phi\,$ = surface admitting a light sheet closest to bdy



- Maximin surface [Wall, '12]
 - maximize over minimal-area surface on a spacelike slice
 - requires the entire collection of slices & surfaces
 - implements homology constraint automatically
 - useful for proofs (e.g. SSA)

Covariant re-formulations of HEE

All of these are the same geometrical construct.



BUT it does not elucidate the relation to quantum information:

- Where does the information live?
- Mutual information I(A:B) = S(A) + S(B) S(AB) is given by surfaces located in different spacetime regions.
- Geometric proof of SSA ($S(AB) + S(BC) \ge S(B) + S(ABC)$) obscures its meaning as monotonicity under inclusion of correlations

Bit thread picture of (static) EE

- Reformulate EE in terms of flux of flow lines [Freedman & Headrick, '16]
 - let v be a vector field satisfying $\,
 abla \cdot v = 0 \,$ and $\, |v| \leq 1$. Then EE is given by

$$S_{\mathcal{A}} = \max_{v} \int_{\mathcal{A}} v$$

- By Max Flow Min Cut theorem, equivalent to RT: (bottleneck for flow = minimal surface)
- Useful reformulation of holographic EE
 - behaves more naturally
 - is more computationally efficient
 - ties better to QI quantities
 - provides more intuition
- How does this extend to time-dependent settings?



Natural possibilities



Convex optimization as a tool

- Max-flow/min-cut is an example of Lagrangian duality in theory of convex optimization
- Setup:
 - Convex program P: minimize $f_0(y)$ over $y \in \mathcal{D}$ such that $\forall i, f_i(y) \leq 0, \forall j, h_j(y) = 0$

convex domain

convex functions

affine functions

• More general problems may be converted to the requisite form via convex relaxation

- use Lagrange multipliers $L(y, \lambda, \nu) \equiv f_0(y) + \sum_i \lambda_i f_i(y) + \sum_i \nu_j h_j(y)$
- solution via convex optimization: $p^* = \inf_y \sup_{\lambda,\nu} L(y,\lambda,\nu)$
- Lagrangian duality: swap order
 - new extremization problem, in new variables

Convex optimization as a tool

- Strategy:
 - Formulate the (Lorentzian) min cut side as convex relaxation
 - Interpret the dual geometrically
- Eg. for static case:



Flow lines

 Via HRT → maximin → maximax → convex relaxation & Lagrangian duality, we still get flow lines = covariant bit threads...



But what is the QI interpretation ?

- Entanglement entropy counted by events ?
 - e.g. # of indep. measurements that can be performed within D[A]
 - novel interpretation...
- Why are I-d structures natural?
 - why is a specific measurement connected to another instantaneous event somewhere in \mathcal{A}^c ?

Summary & Outlook

- Holography conveniently geometrizes entanglement
 - Finding bulk geometrical constructs is (relatively) easy!
 - Useful in proving important properties!
 - Why is EE related to geometry so simply?
 - Duals of other measures of entanglement?
- General covariance is a powerful guiding principle
 - Motivated subregion/subregion duality
 - Covariantize bit threads to elucidate essence of holographic EE
 - Significance of instantaneous nature: (Why) I -d threads?
- Convex relaxation and Lagrangian duality is a powerful tool
 - Motivates new geometric constructs, new elegant proofs, connections...
 - Other applications?
- Relation between spacetime (gravity) and entanglement?

