

Theoretical *Challenges* in the Post-Planck Era

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Bay Area Particle Theory Seminar

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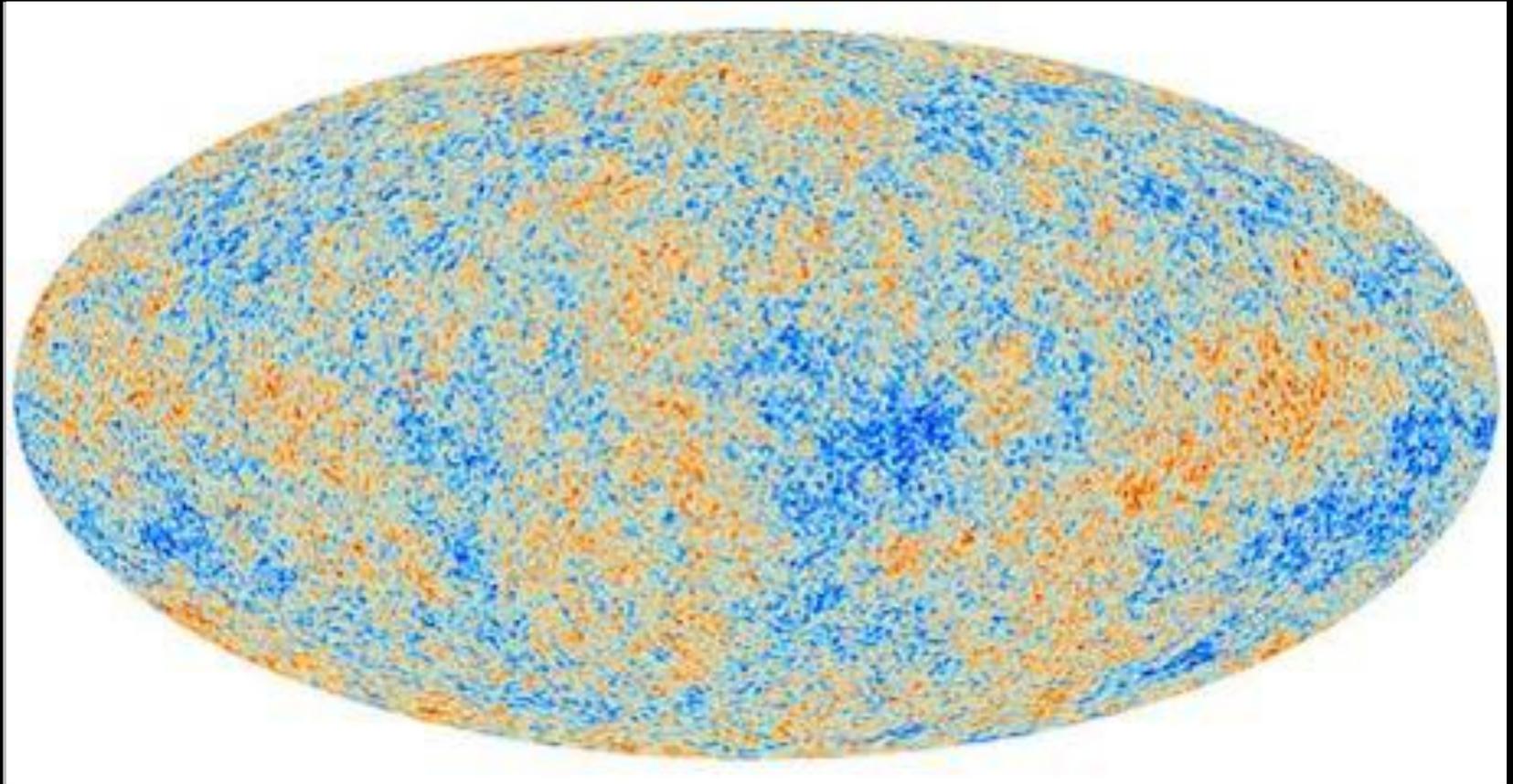
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- I would like to challenge you to consider an unpopular idea designed to address one of these open problems

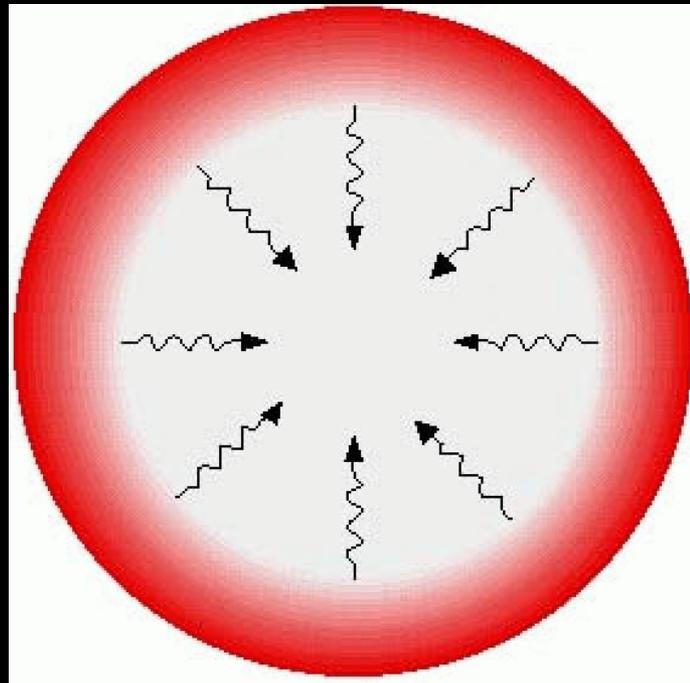
Theoretical *Challenges* in the Post-Planck Era

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- Are we living in a boring Universe?
- I would like to challenge you to consider an unpopular idea designed to address one of these open problems
- Please challenge me if you think this idea is unpopular for a good reason!

“What are we to make of the Planck results?” (MP)



We see photons today from last scattering surface at $z=1100$



Acoustic Oscillations

Pressure of radiation acts against clumping

If a region gets overdense, pressure acts to reduce the density: **restoring force**



As a result of the pressure of radiation

Pressure of radiation

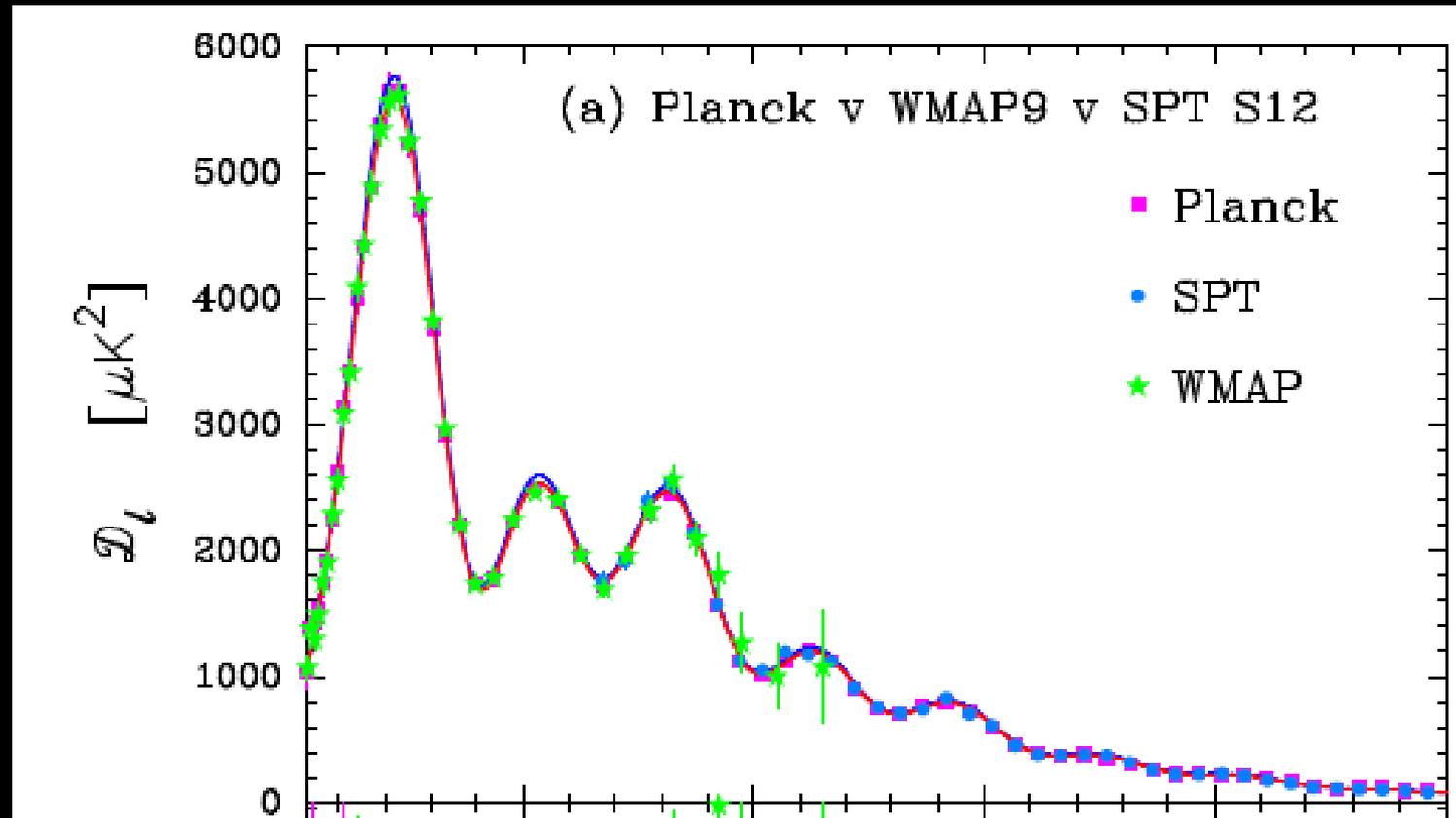
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If a region gets over
density: restoring f

to reduce the



Hence peaks and troughs in Spectrum: exactly like harmonics associated with musical instruments



CMB is different because ...

- Fourier Transform of spatial, not temporal, signal

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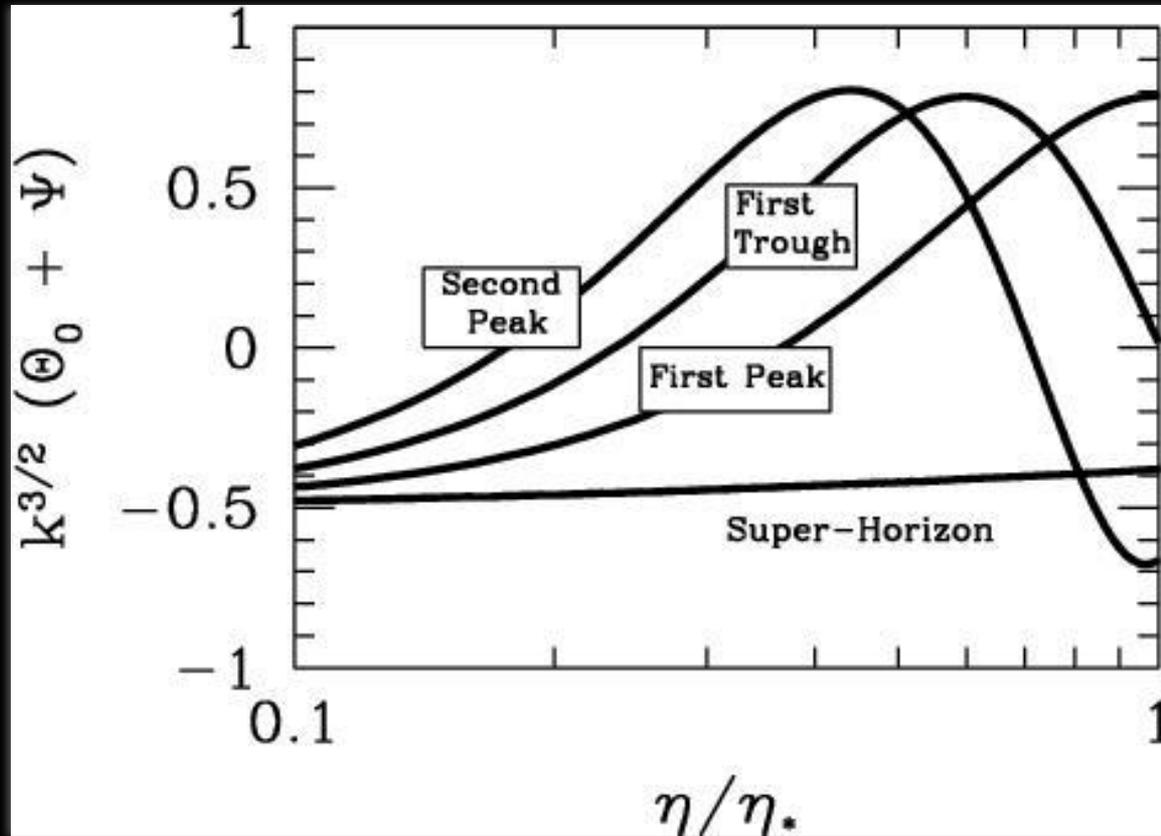
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- Time scale much longer (400,000 years vs. 1/1000 seconds)

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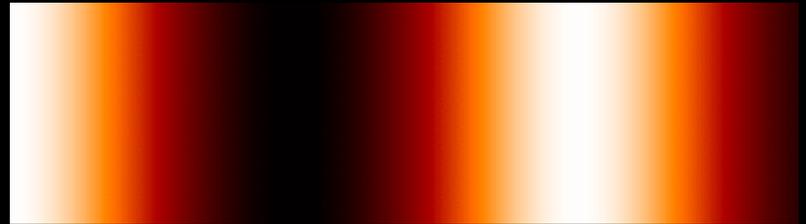
- Fourier Transform of spatial, not temporal, signal
- Time scale much longer (400,000 years vs. 1/1000 seconds)
- No finite length: all frequencies allowed!

Why peaks and troughs?

- **Vibrating String:** Characteristic frequencies because ends are tied down
- **Temperature in the Universe:** Small scale modes enter the horizon earlier than large scale modes

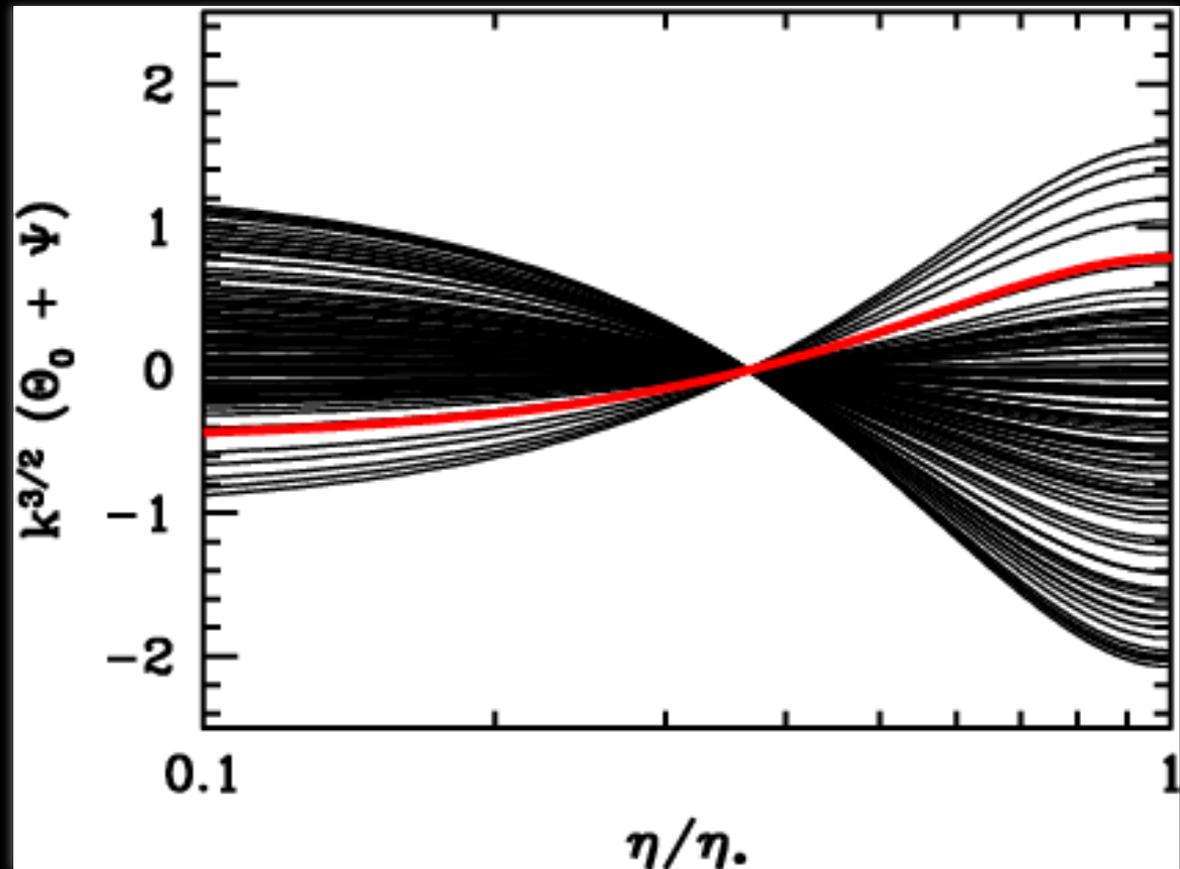


BUT: there are an infinite number of modes associated with a given wavelength. The CMB first peak, first example, comes from a sum over an infinite number of Fourier modes, each with a different orientation.



Interference should destroy peak structure

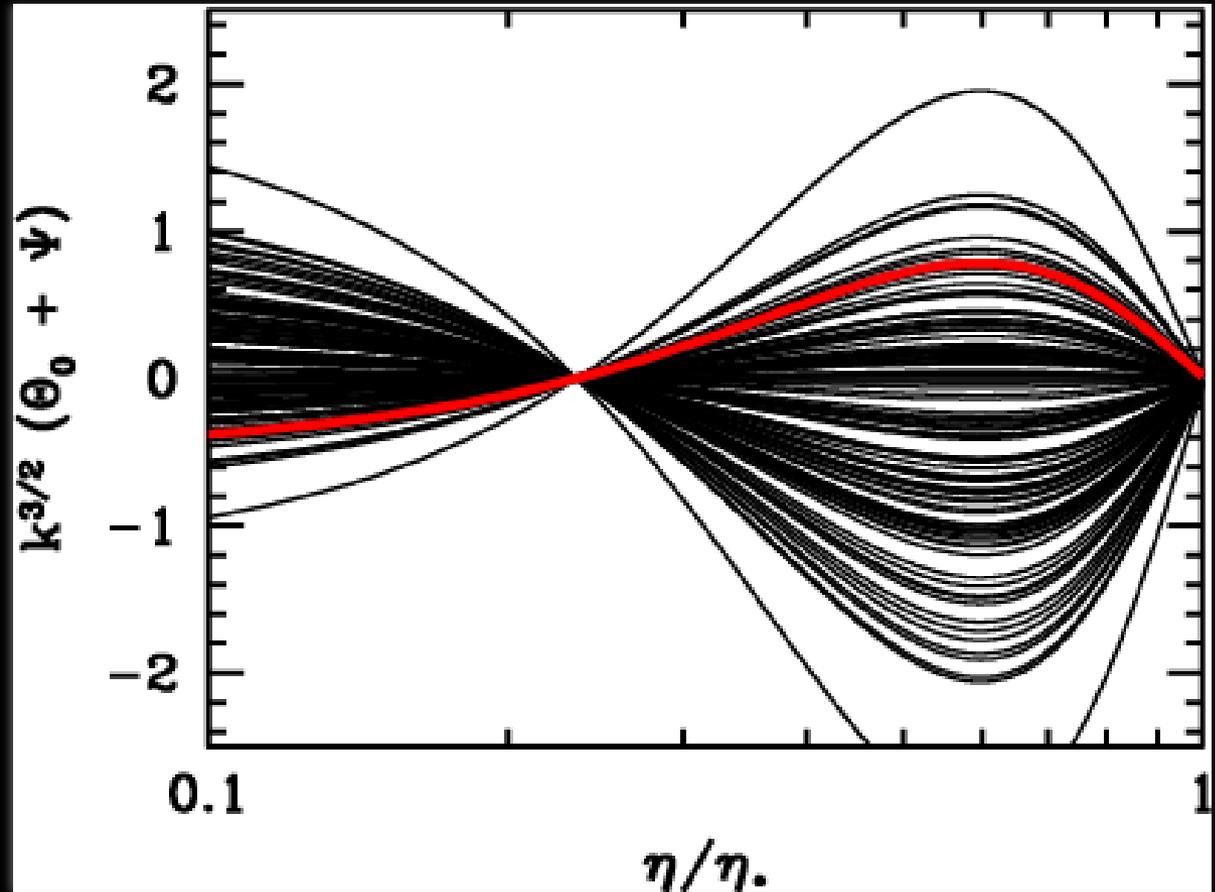
There are many, many modes with similar values of k . All have different initial amplitude. Why are all in phase?



First Peak Modes

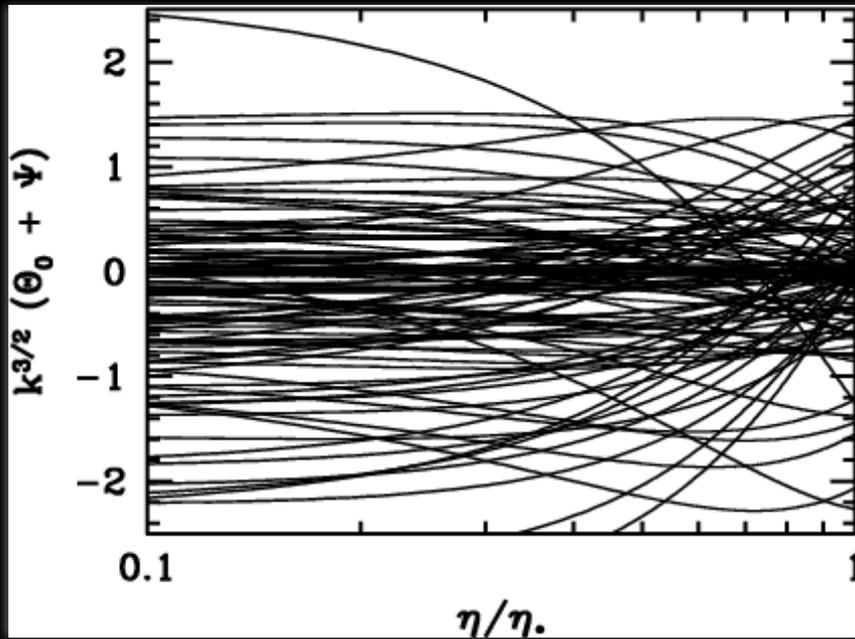
An infinite number of violins are synchronized

Similarly, all modes corresponding to first trough are in phase: they all have zero amplitude at recombination. Why?

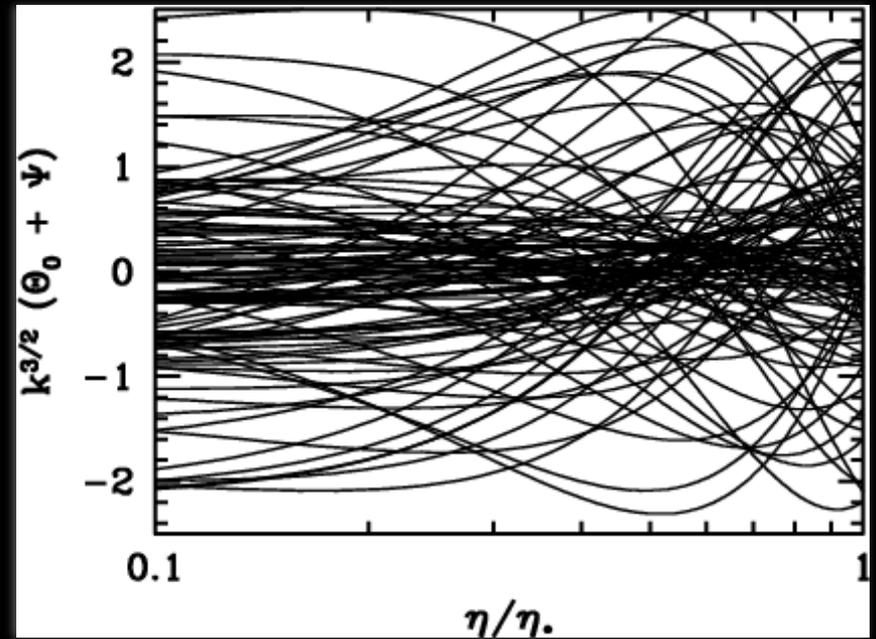


Without synchronization:

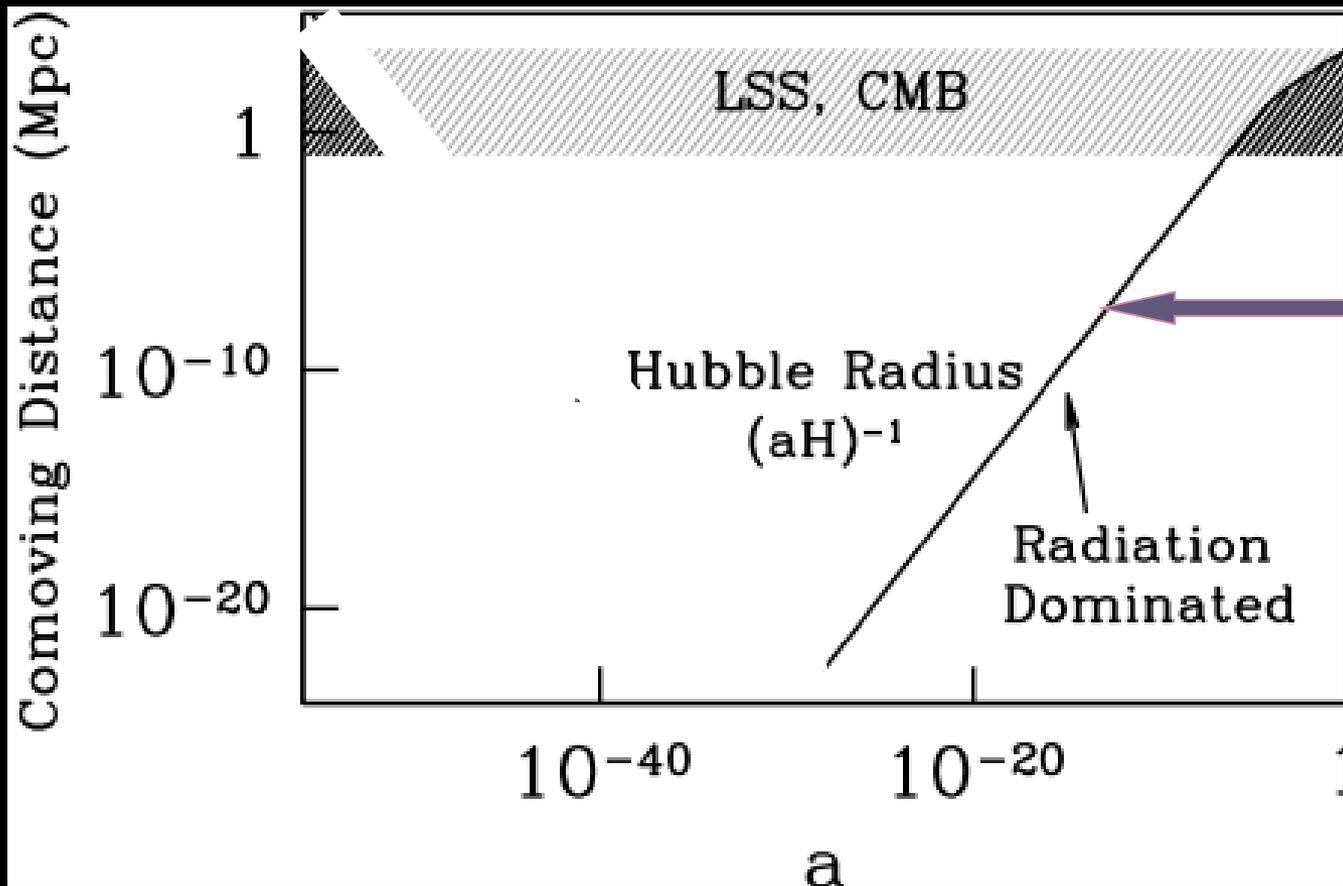
First "Peak"



First "Trough"

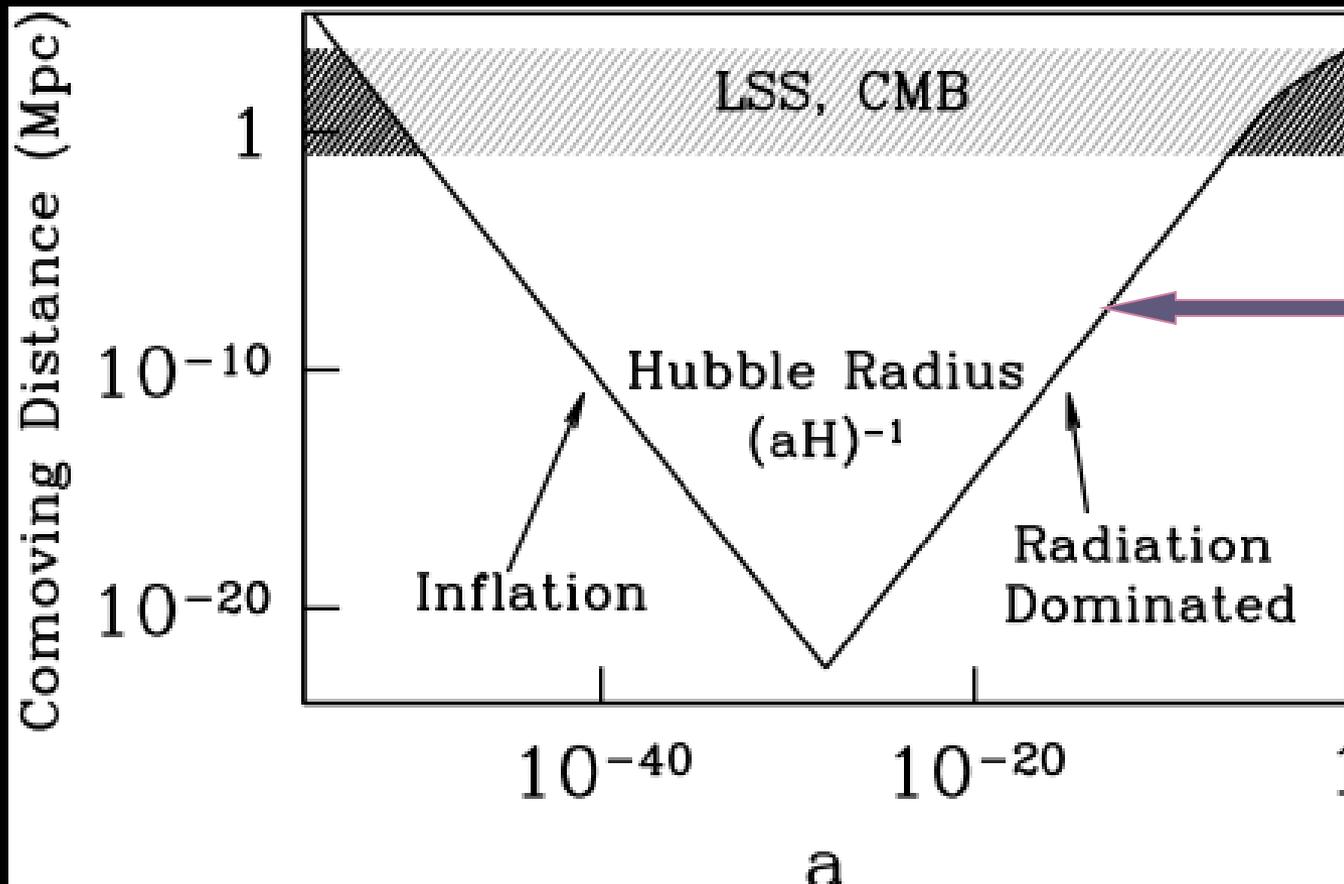


It is worse than this: The Hubble radius is the distance over which information can propagate causally at a given epoch



Perturbations
*outside the
horizon*

Inflation is an epoch during which the comoving Hubble radius decreases: perturbations were synchronized before/during inflation



Perturbations
*outside the
horizon*

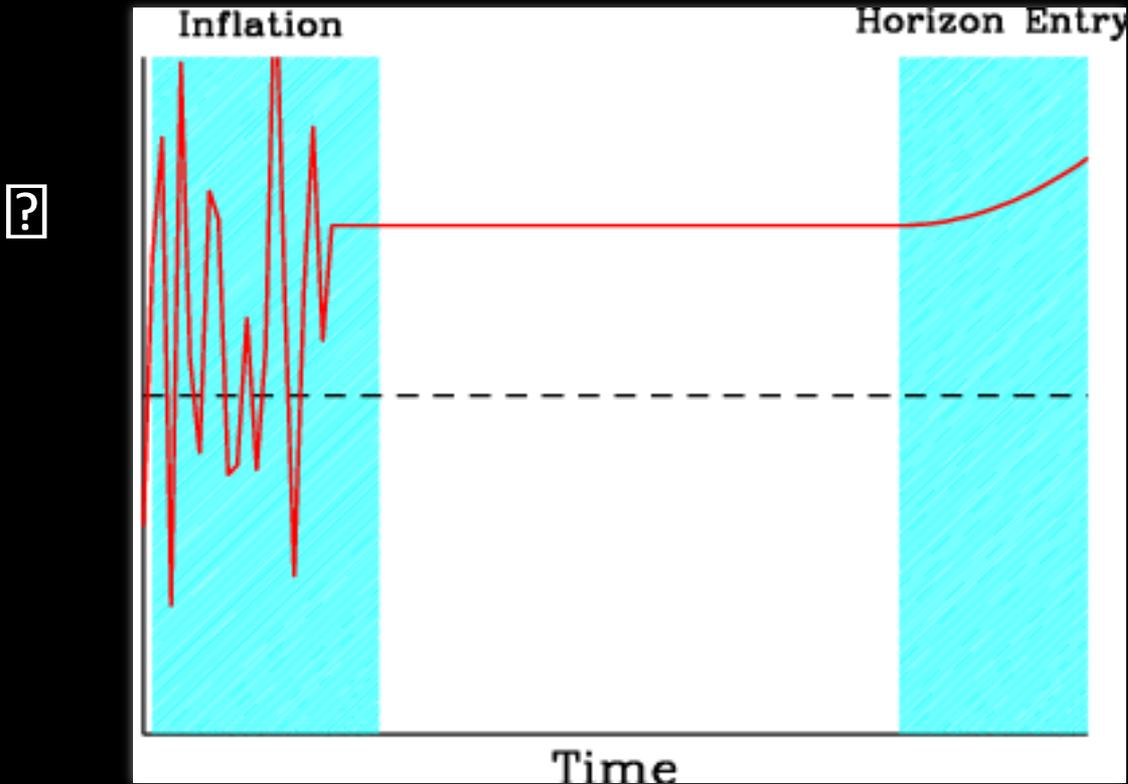
Inflation produces perturbations

- Quantum mechanical fluctuations

$$\langle \Psi(\mathbf{k}) \Psi(\mathbf{k}') \rangle = \frac{1}{(2\pi)^3} \delta^3(\mathbf{k} + \mathbf{k}') P_{\Psi}(\mathbf{k})$$

- Inflation stretches wavelength beyond horizon: $\Psi(\mathbf{k}, t)$ becomes constant
- Infinite number of independent perturbations w/ independent amplitudes

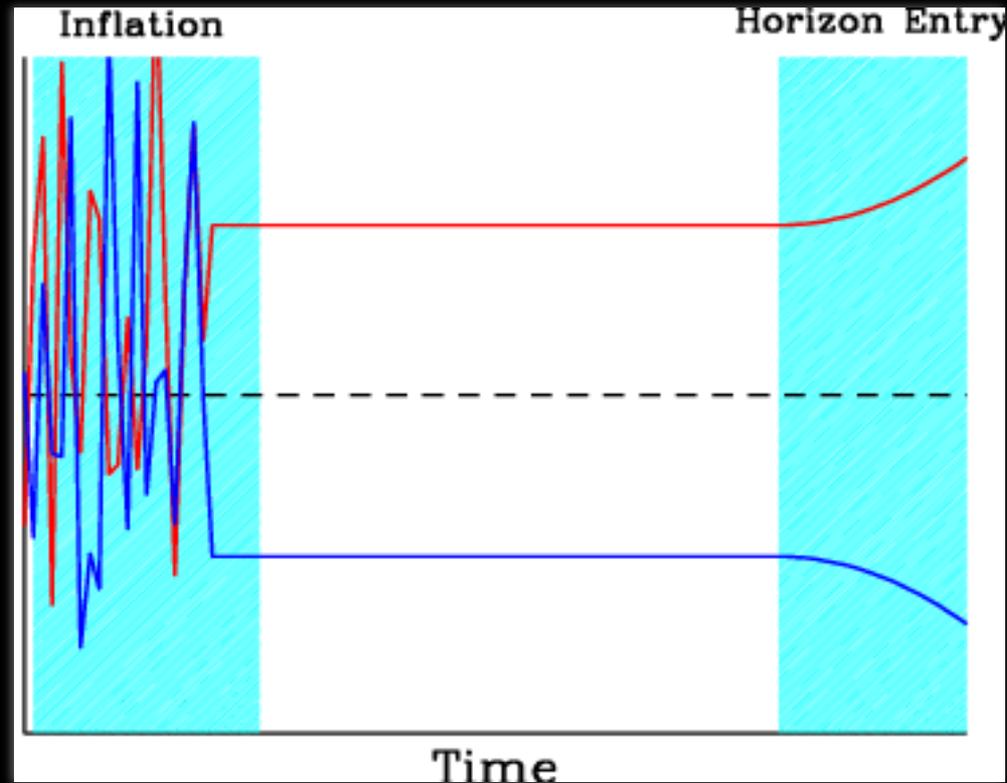
Inflation synchronizes all modes



All modes remain constant until they re-enter horizon.

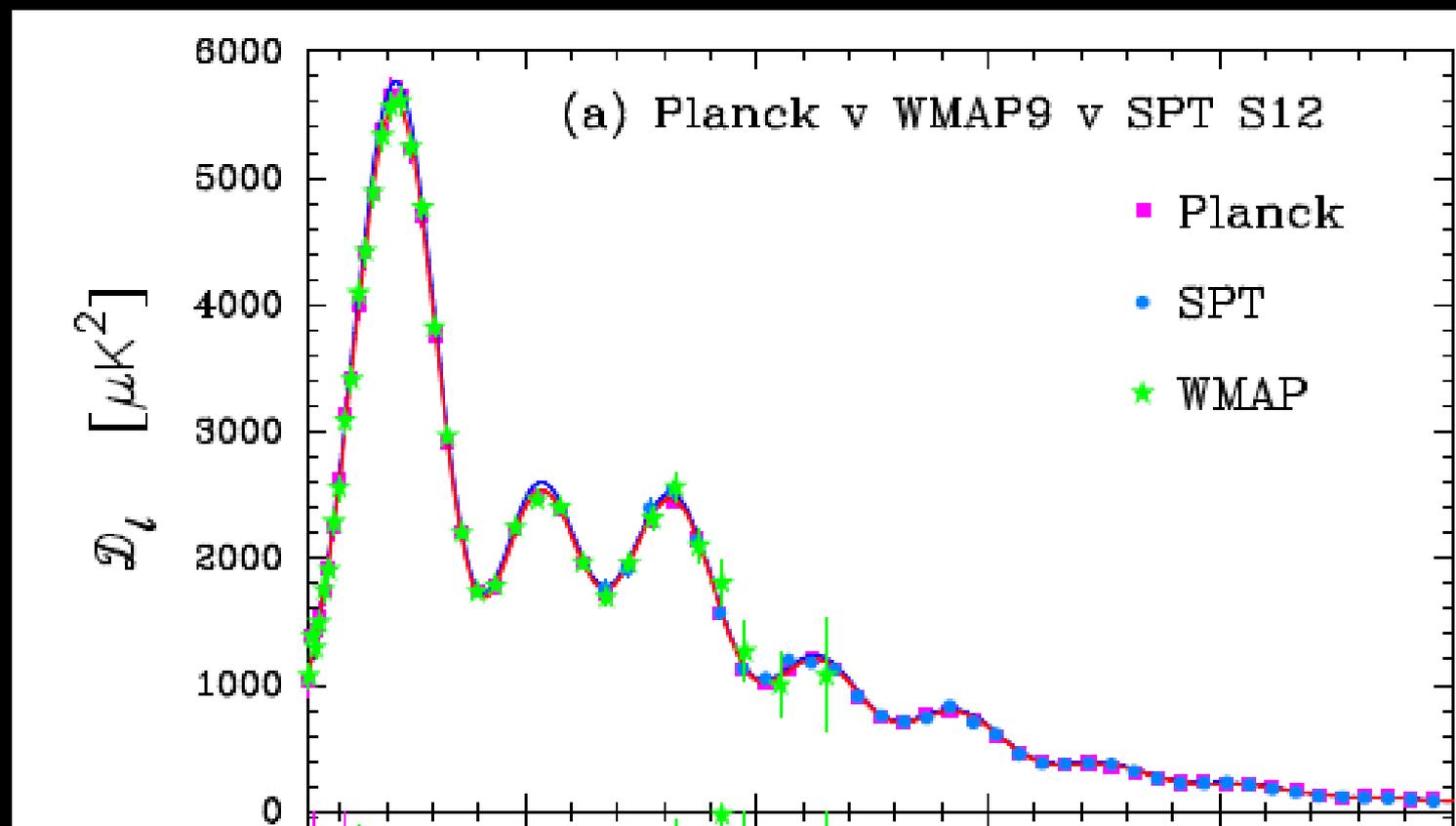
Inflation synchronizes all modes

?



All modes remain constant until they re-enter horizon.

Coherence of Peaks and Troughs Strong Evidence for Inflation (or something even crazier)



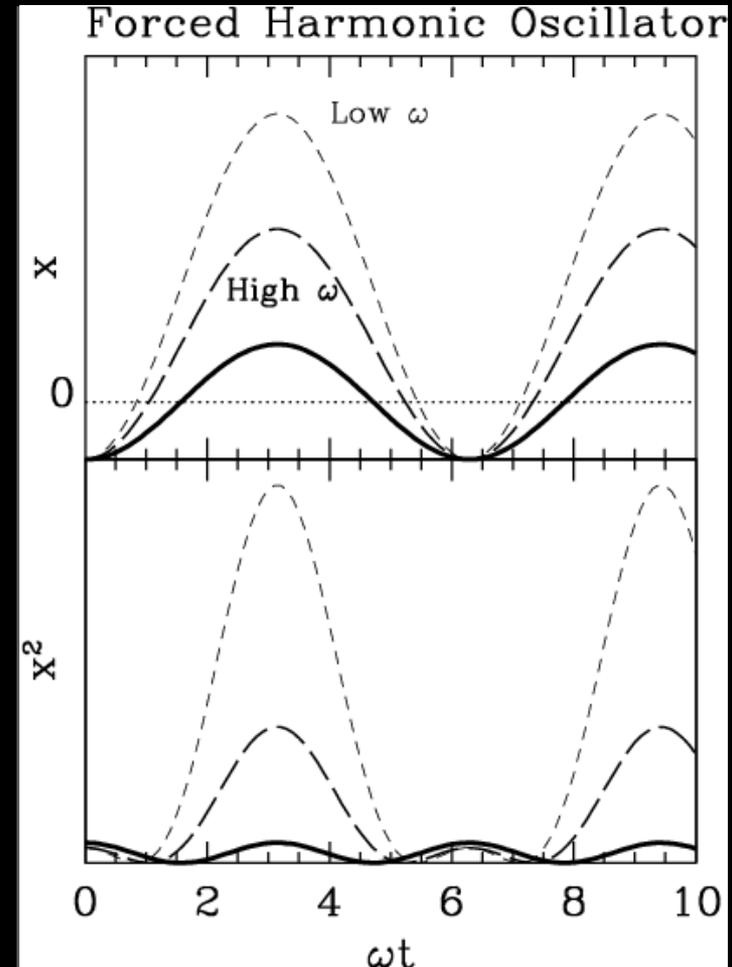
Harmonics carry information about the “instrument”

$$\ddot{x} + W^2 x = F$$

with frequency inversely proportional
to baryon density

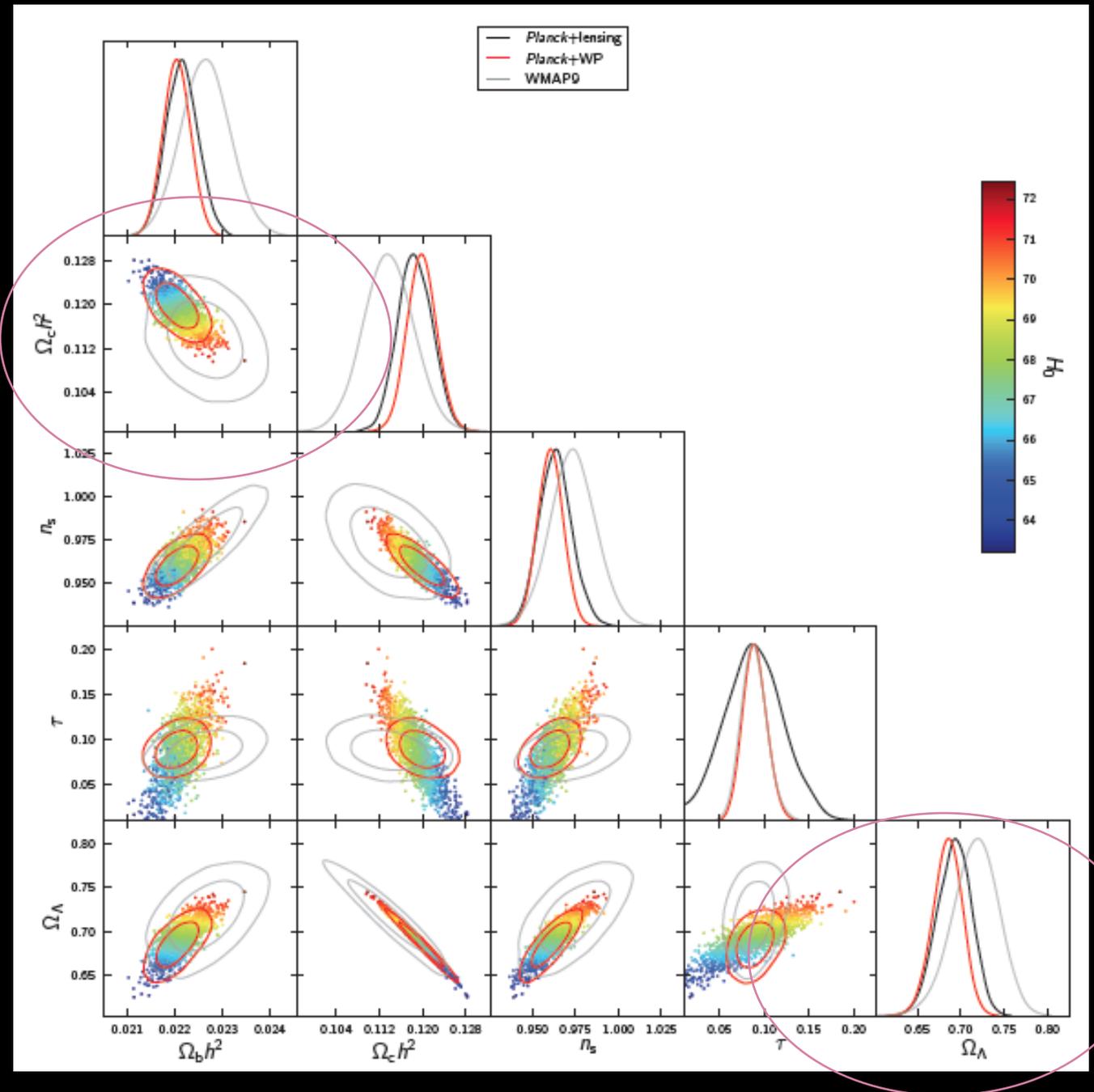
$$\omega = kc_s = \frac{k}{\sqrt{3(1 + 3\rho_b / 4\rho_\gamma)}}$$

Higher ω (e.g. fewer baryons) \rightarrow no
odd/even peak disparity.



Robust evidence for non-baryonic dark matter

Independent evidence for dark energy



Challenge: What is this new physics?

- ***Non-Baryonic Dark Matter***: Popular framework with many free parameters. If detected, lots of stuff to do. If not, when do we give up on SUSY and WIMP's? What replaces them?
- ***Inflation***: Handful of models have risen to prominence. Disappointment that PNG is not large. Wait for B-modes. Then what?
- ***Dark Energy***

Modified Gravity to explain Accelerating Universe

Can explain acceleration without dark energy (w) by modifying GR:

$$S = \frac{1}{16\pi G} \int d^4x \sqrt{-g} [R + f(R)] + \int d^4x \sqrt{-g} L_m$$

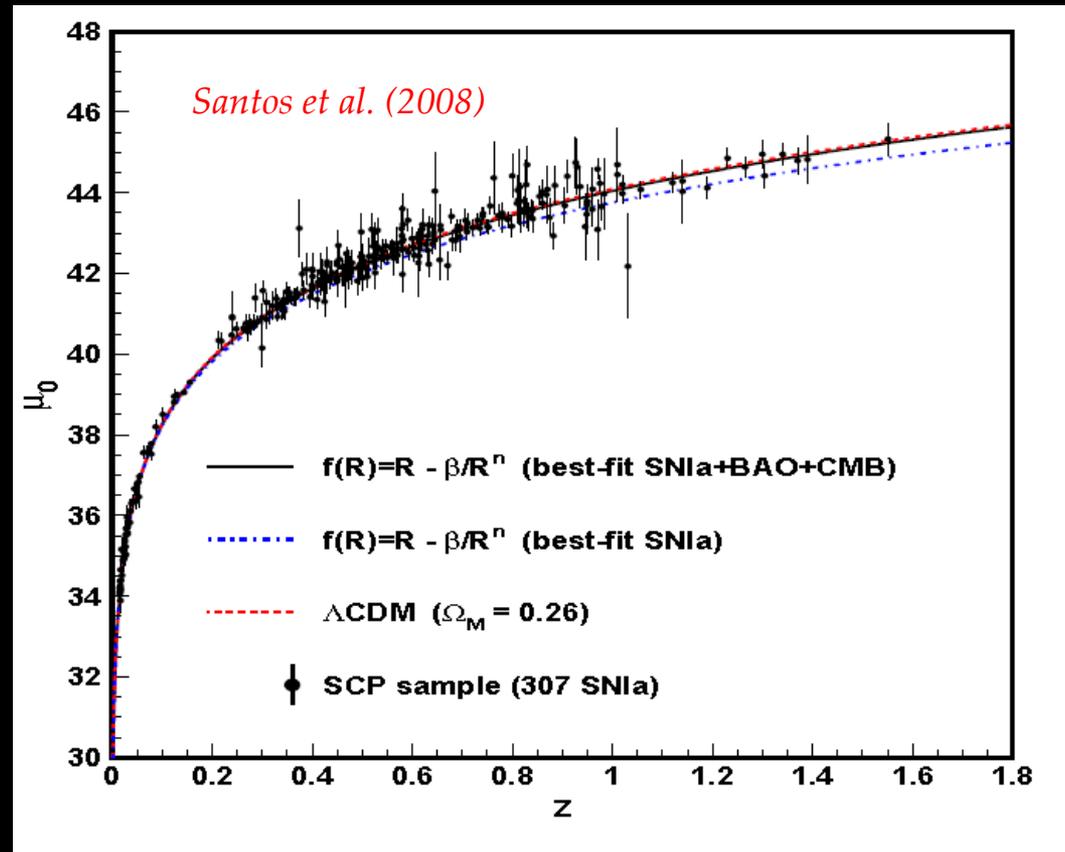
For the cosmological metric, the acceleration equation generalizes to:

Get acceleration if these terms are positive

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} (\rho + 3P) + \left[\frac{\partial f}{\partial R} H^2 - \frac{f}{6} - \frac{\partial^2 f / \partial R^2}{2} \right]$$

Easy to fit Supernova Data

More generally, MG models can fit any redshift-distance relation measured by SN and other cosmological probes



$f(R)$ has dimensionful parameter order 10^{-33} eV

Scalar Tensor Models

- $f(R)$ models are a subset of *scalar-tensor* models
- Challenge the implicit assumption of General Relativity that the metric in the Einstein-Hilbert action is the same as the metric that couples to matter

$$S_{EH} = \frac{1}{16\pi G} \int d^4x \sqrt{-\tilde{g}} R[\tilde{g}] \quad S_m = \int d^4x \sqrt{-g} L_m$$

- Allow $\tilde{g} = e^{-2\phi/m_p} g$
- Scalar-Tensor model is defined by dynamics $S[\Phi]$, related to $f(R)$
- Extra degree of freedom can cause problems in Solar System: require *screening* mechanism to suppress in high-density/small scale regions

Massive graviton

- Fierz & Pauli (1939) first introduced massive graviton; recent interest as cause of acceleration
- Two new degrees of freedom \leftrightarrow Two problems:
 - Solar System constraints
 - Ghosts (negative kinetic energy term leads to instabilities)
- Nonlinear Fierz-Pauli action produces a screening mechanism
- Subset of nonlinear terms in Fierz-Pauli action ghost-free:
Galileons

Distinguishing General Relativity + Dark Energy from Modified Gravity

- Fix expansion history in both models (Zero order cosmology)
- Compare how structure in the universe grows (Perturbations)

Perturbations in Modified Gravity

Lue, Scoccimarro, and Starkman (2004); Bertschinger (2006); Hu & Sawicki (2007)

Start with the perturbed FRW metric

$$ds^2 = -(1 + 2\Psi)dt^2 + a^2(t)(1 - 2\Phi)(dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2)$$

Generally two differences between MG and GR:

GENERAL RELATIVITY

MODIFIED GRAVITY

$$\Phi - \Psi = 0$$

$$\Phi - \Psi \neq 0$$

$$\nabla^2 \Phi = -4\pi G \rho_m a^2 \delta$$

$$\nabla^2 \Phi = -4\pi G_{\text{eff}} \rho_m a^2 \delta$$

Deviations from GR

Closer to observations to
define:

$$\begin{aligned}\Psi &\equiv [1 + \mu(k, a)] \Psi_{GR} \\ \Psi - \Phi &= [1 + \Sigma(k, a)] [\Psi_{GR} - \Phi_{GR}]\end{aligned}$$

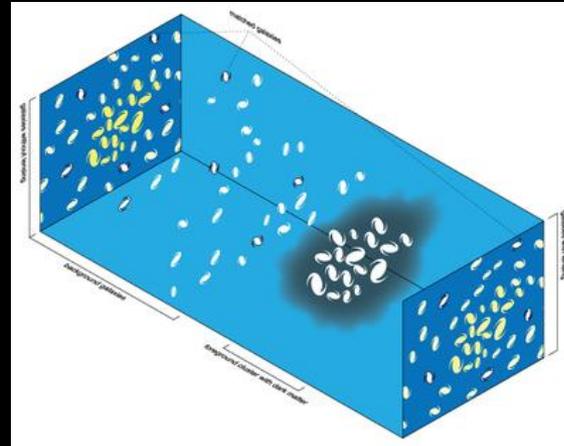
Massive particles
respond to μ , while
massless particles
(photons) respond to Σ

Deviations from GR

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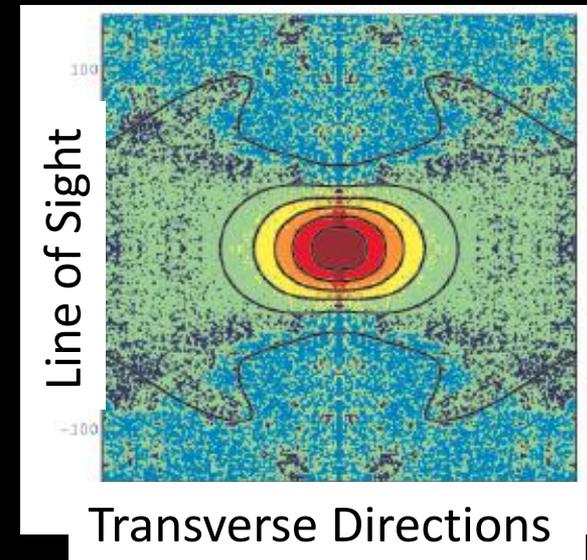
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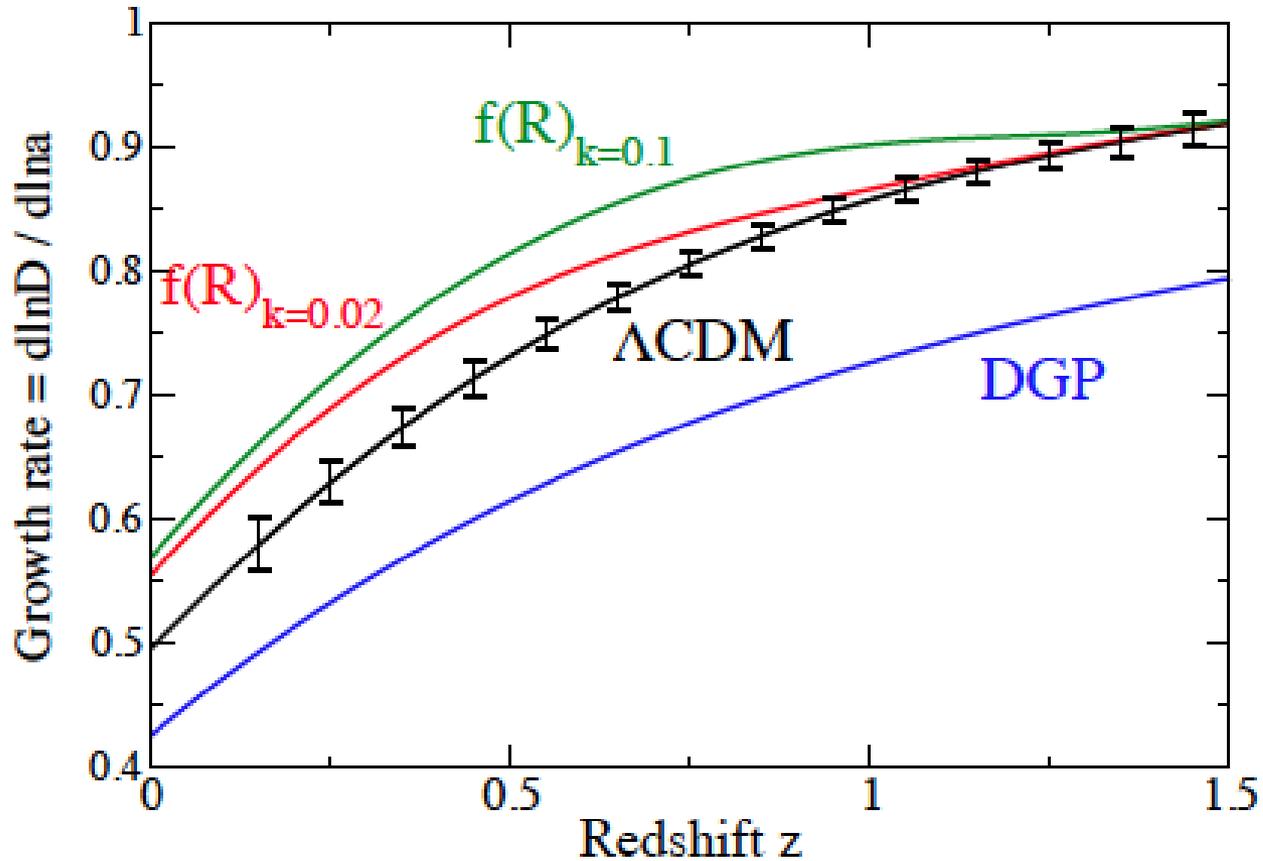


Lensing can measure Σ

Redshift Space Distortions sensitive to μ

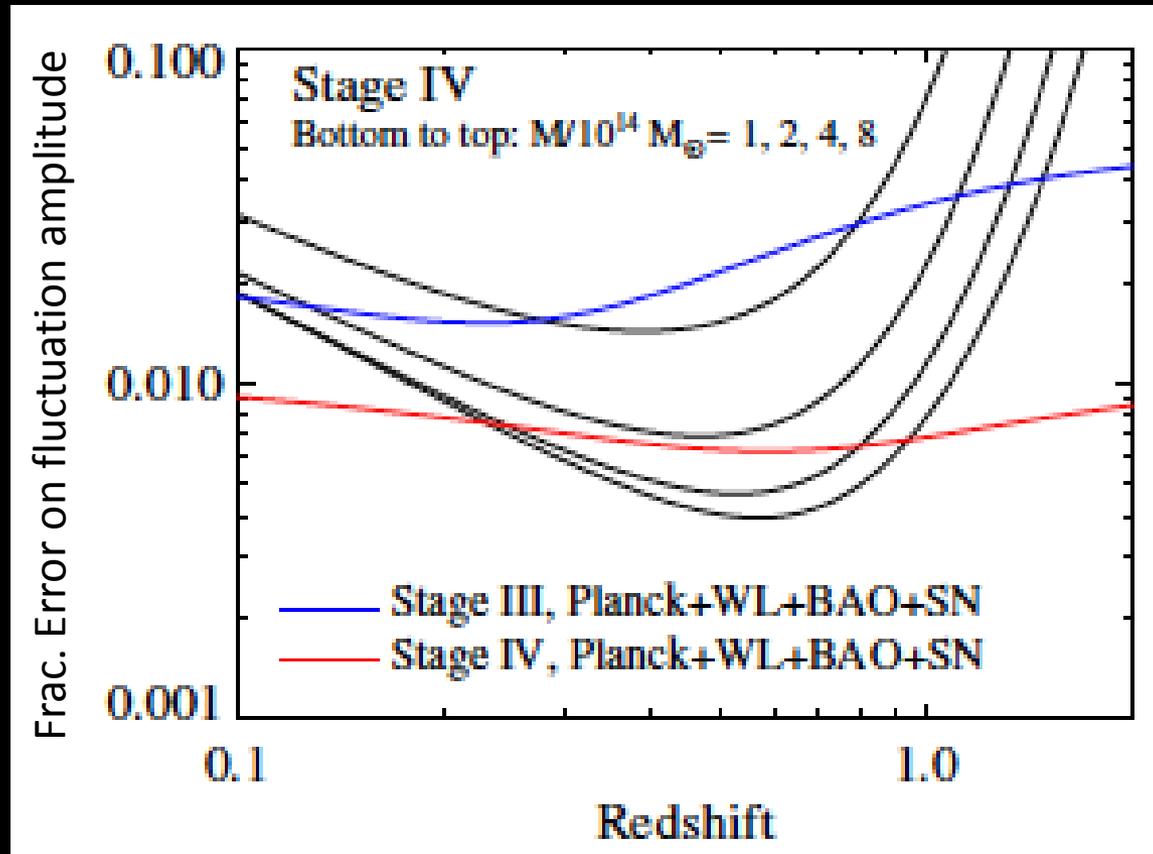


Projected constraints from RSD



Huterer et al. 2013

Projected Constraints from Lensing



Example: Nonlocal Gravity

(Deser and Woodard 2007)

$$\mathcal{L}_g = \frac{1}{16\pi G} \sqrt{-g} R \left[1 + f\left(\frac{1}{\square} R\right) \right]$$

- Argument of f is dimensionless: no new parameters (almost all DE/MG models have a mass parameter of order 10^{-33} eV)
- Deviations occur only at late times, logarithmically growing from the epoch of equality
- Easy to evade Solar System constraints
- Terms like this generated by quantum corrections in string theory away from critical dimension (Polyakov 1981). Banks (1988) has more general arguments
- Free function f can be chosen to fit any expansion history

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- **Free function f**

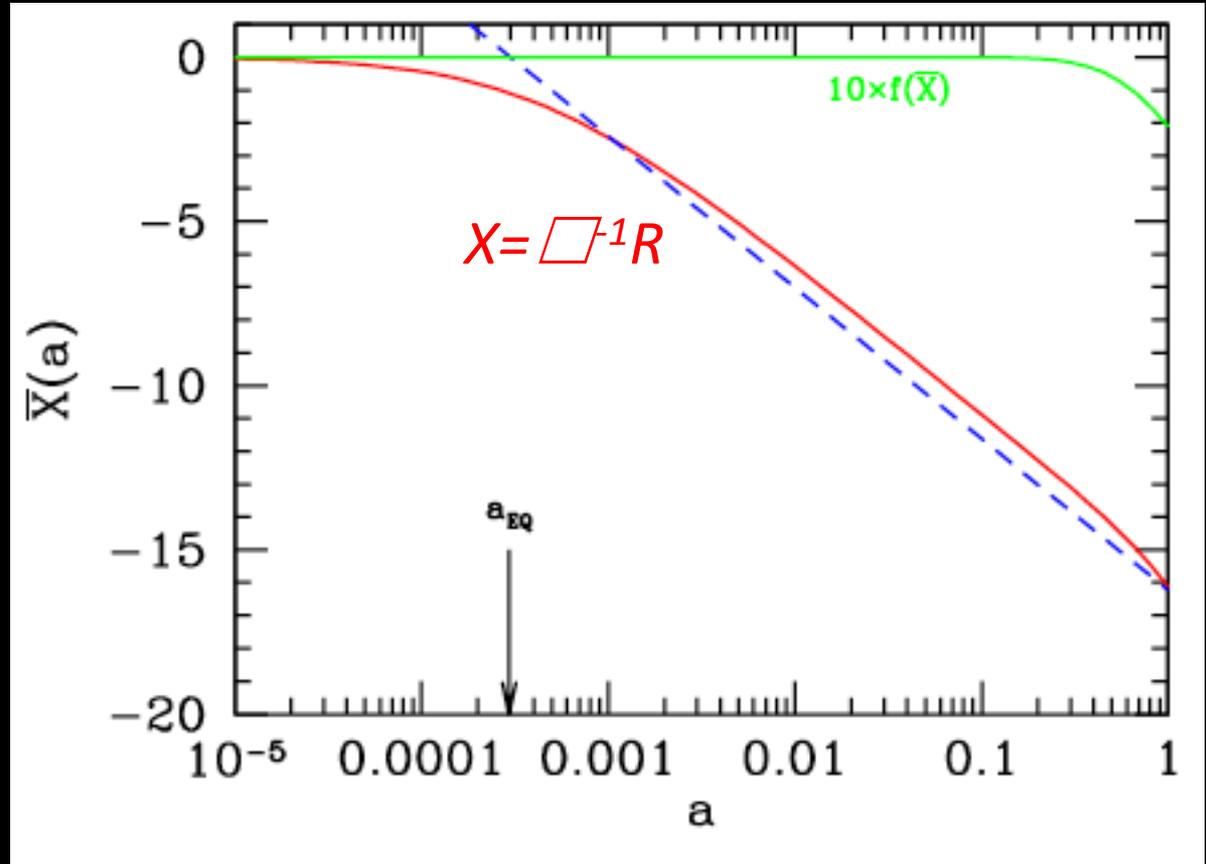
Nonlocal Gravity

Leads to new terms in Einstein's equation

$$\begin{aligned} \Delta G_{\mu\nu} = & \left[G_{\mu\nu} + g_{\mu\nu} \square - D_\mu D_\nu \right] \left\{ f(X) + \frac{1}{\square} [Rf'(X)] \right\} \\ & + \left[\delta_\mu^{(\rho} \delta_\nu^{\sigma)} - \frac{1}{2} g_{\mu\nu} g^{\rho\sigma} \right] \partial_\rho (X) \partial_\sigma \left(\frac{1}{\square} [Rf'(X)] \right). \end{aligned}$$

Nonlocal Gravity

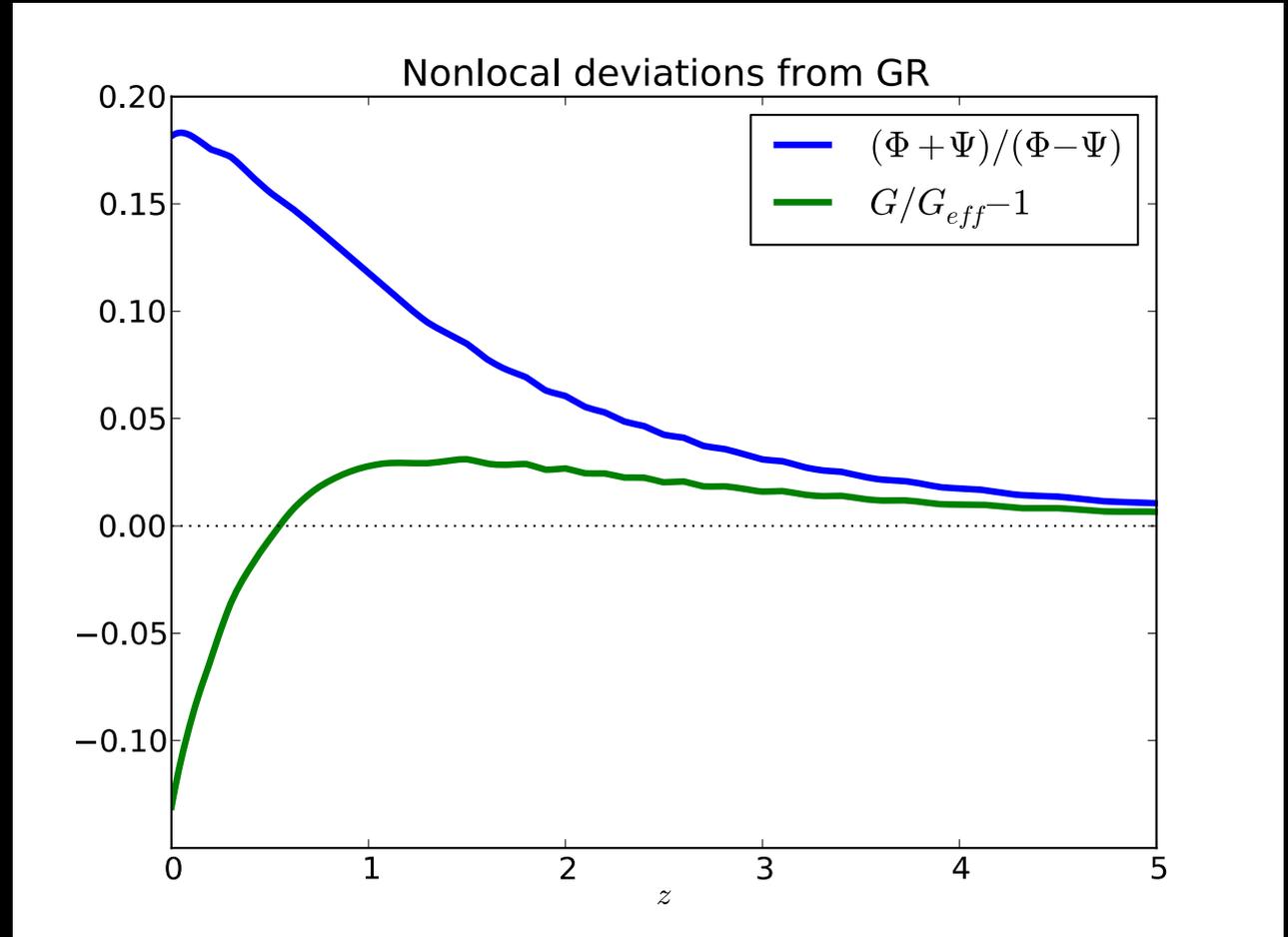
- Solve the new equations in homogeneous cosmology
- Fix free function by matching observed redshift-distance relation (Deffayet & Woodard 2009)



Nonlocal Gravity: Perturbations

(Park & Dodelson 2012; Dodelson & Park 2013)

Gravitational force is weaker at early times but then grows stronger

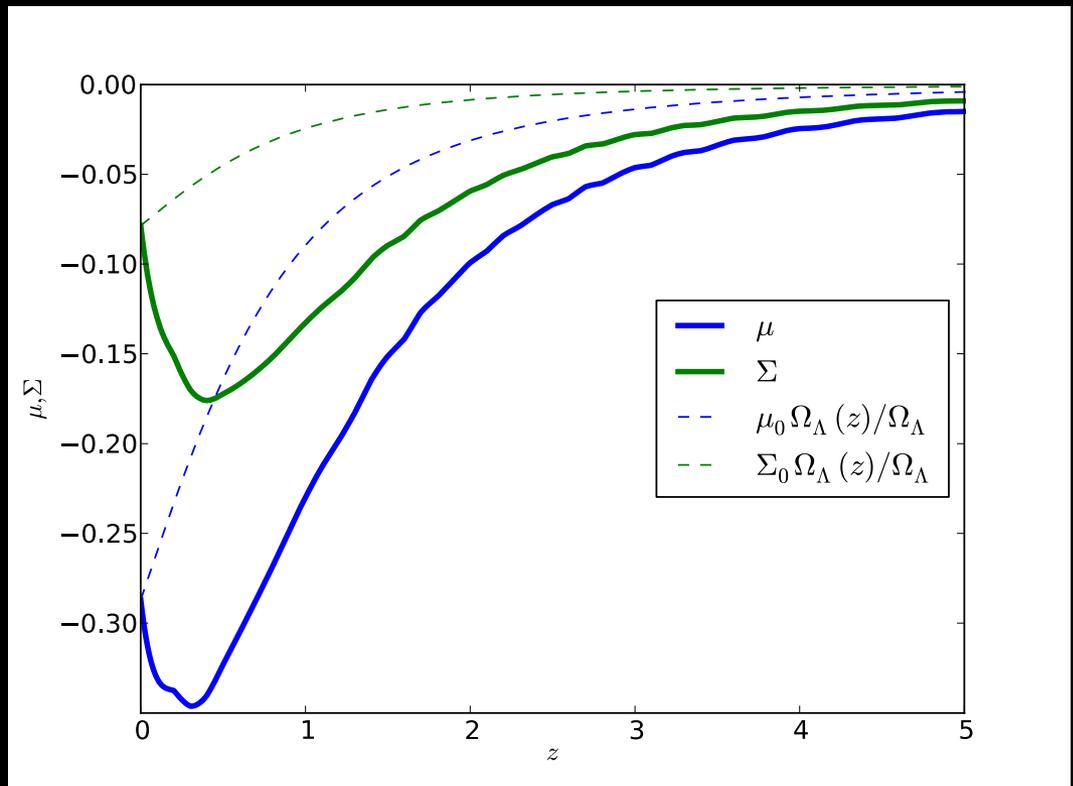


Nonlocal Gravity: Deviations from GR

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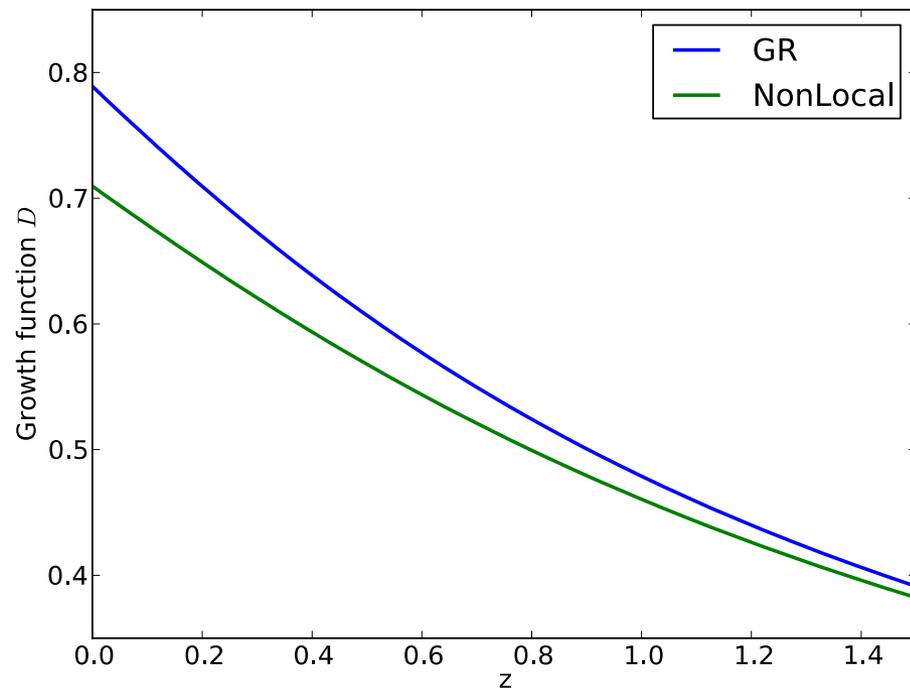
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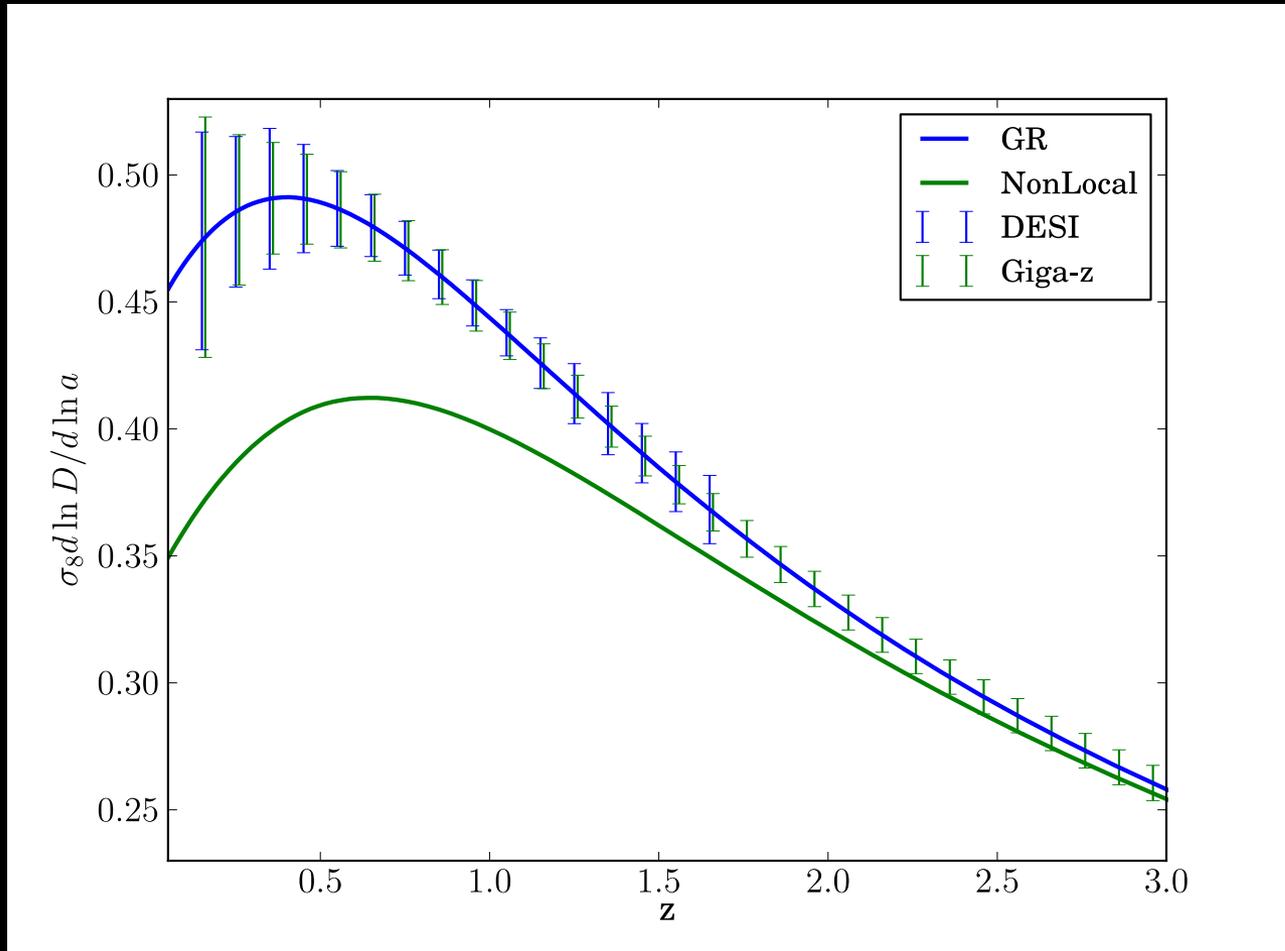
Nonlocal Gravity: Growth Function

$$\frac{d^2\delta}{da^2} + \left[\frac{d\ln(H)}{da} + \frac{3}{a} \right] \frac{d\delta}{da} - \frac{3}{2} [1 + \mu] \frac{\Omega_m}{E^2(a)a^5} \delta = 0$$

Since μ is non-zero and negative, similar to lower matter density, and perturbations grow slower than in GR

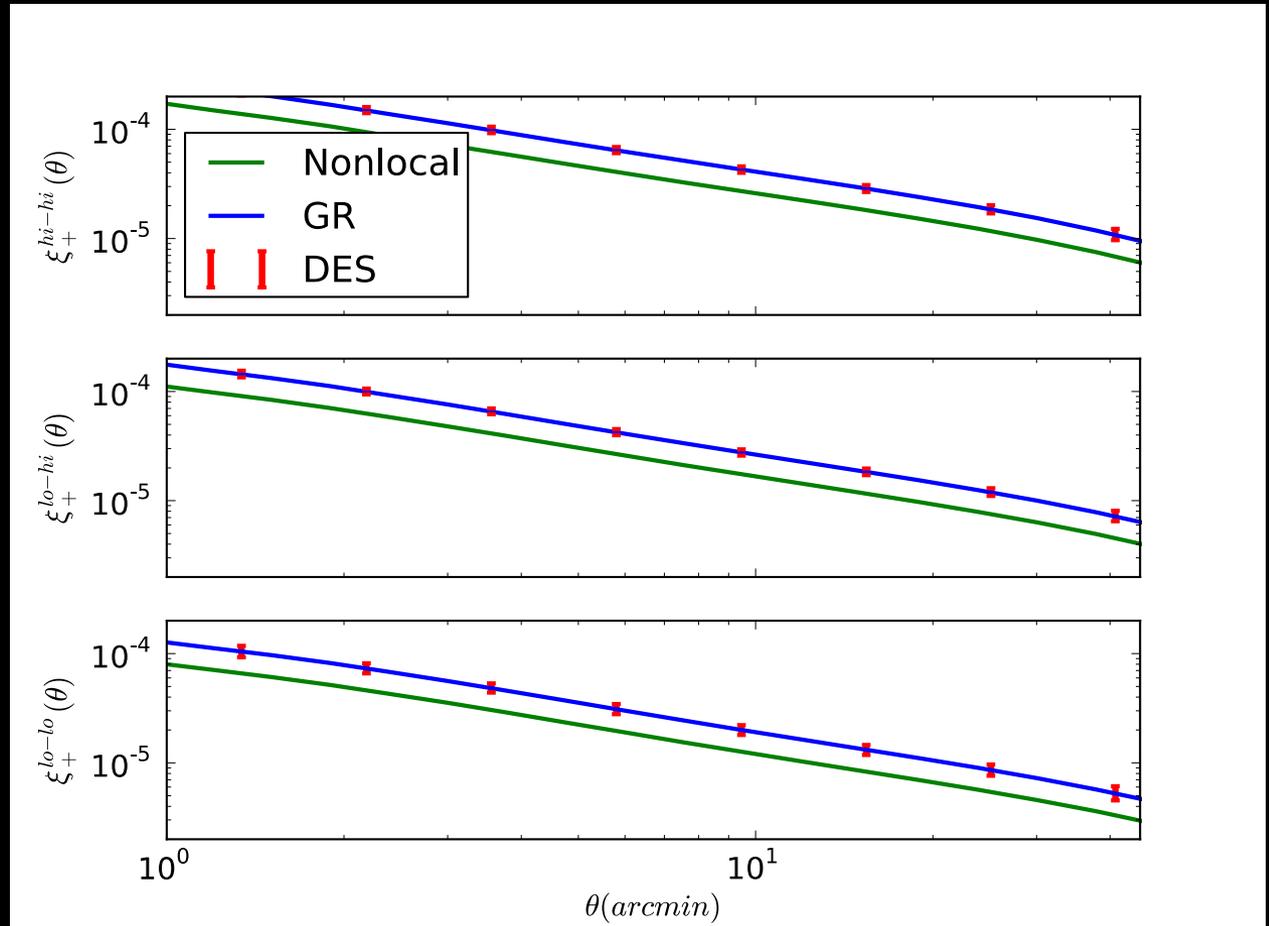


Nonlocal Gravity: Redshift Space Distortions



Nonlocal Gravity: Lensing

Correlation function
of galaxy ellipticities*
in 3 tomographic bins:
Lo-Lo: $z < 0.8$
Hi-Hi: $z > 0.8$
Lo-Hi: Cross



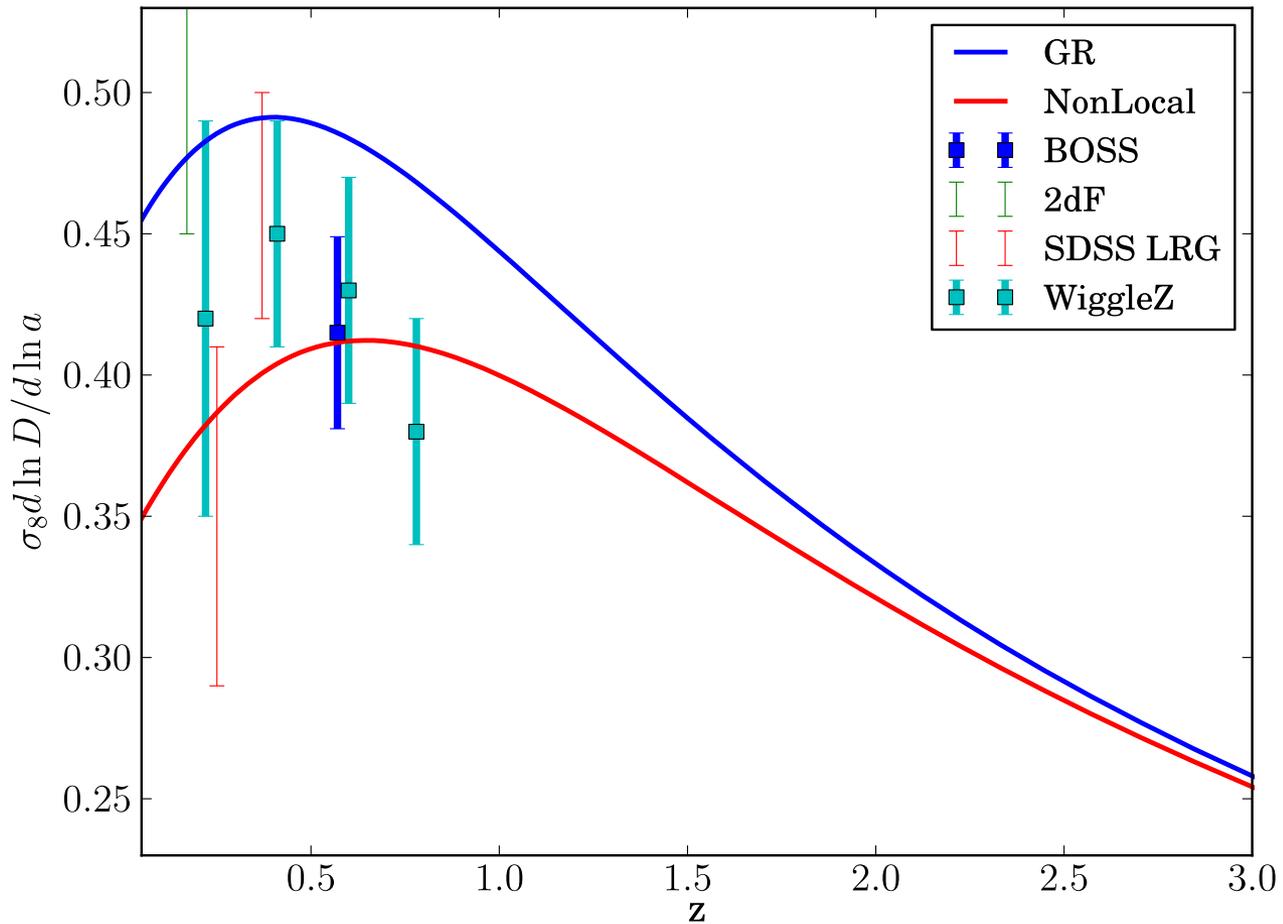
* Distorted by lensing

Just for fun ...

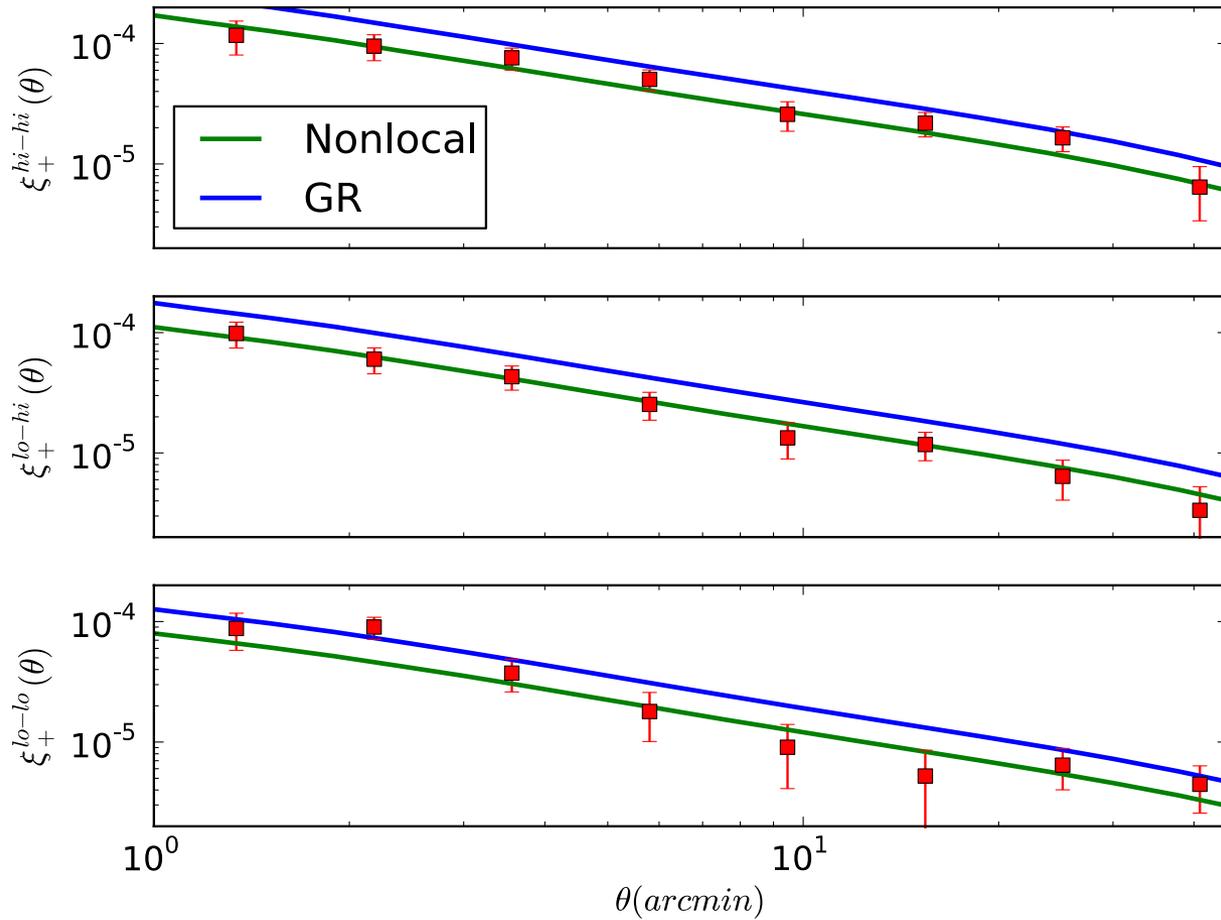
“Stage II” surveys have yielded interesting data already:

- Redshift space distortions (SDSS LRG's, 2dF, Wiggle-Z, BOSS)
- Lensing (SDSS, CFHTLens)

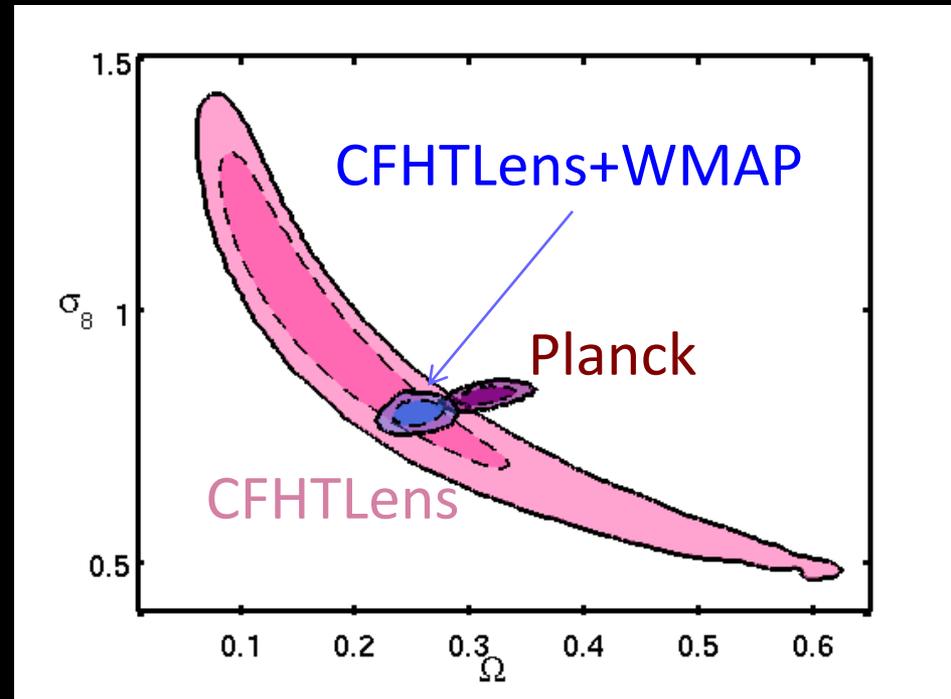
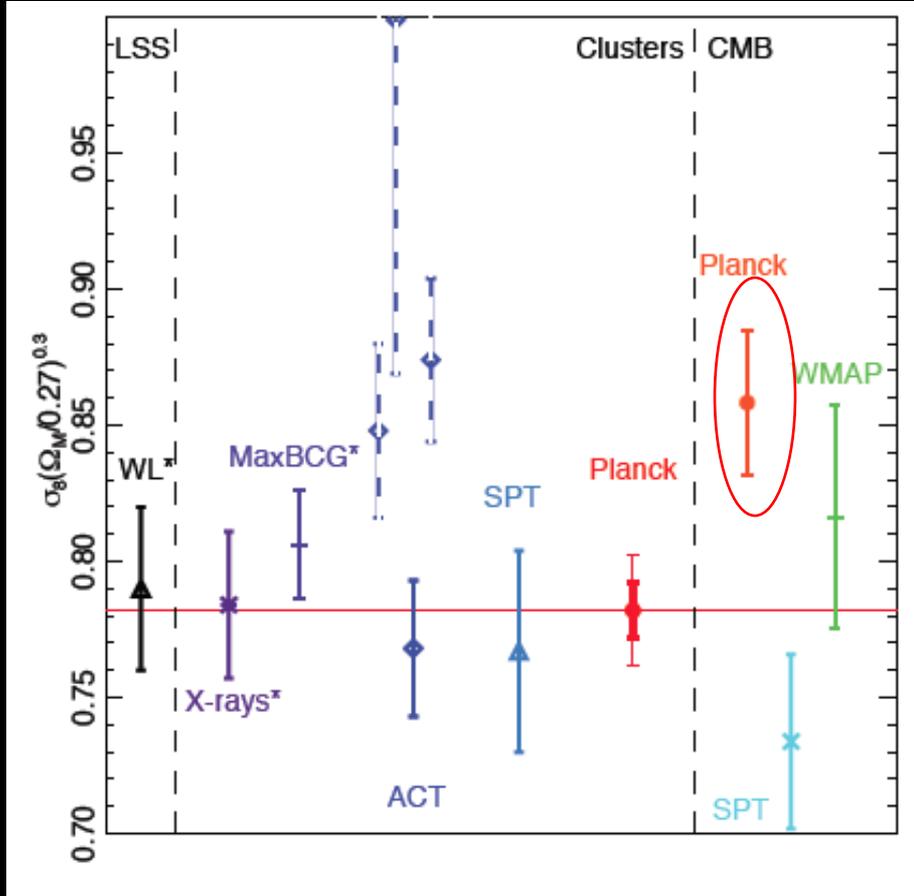
Nonlocal Gravity: Redshift Space Distortions



Nonlocal Gravity: Lensing



Some of this tension has been noted



Planck 2013

Courtesy: Catherine Heymans, CFHTLens

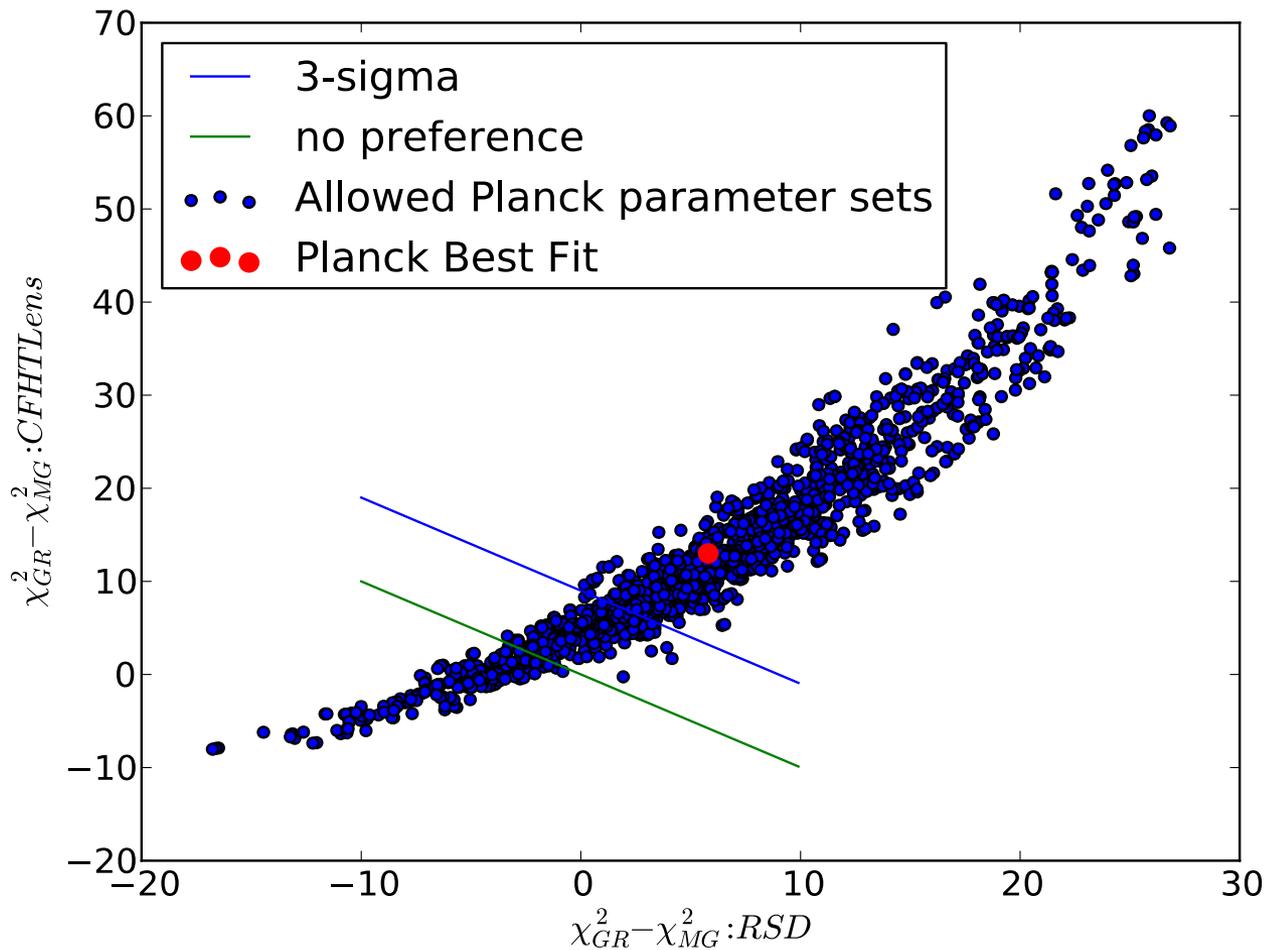
Statistical Significance

With parameters fixed from Planck, no free parameters for either GR or nonlocal gravity. Can compare χ^2 straight up:

RSD points are essentially uncorrelated: $\Delta\chi^2=8.8 \rightarrow$ 3-sigma preference for nonlocal model

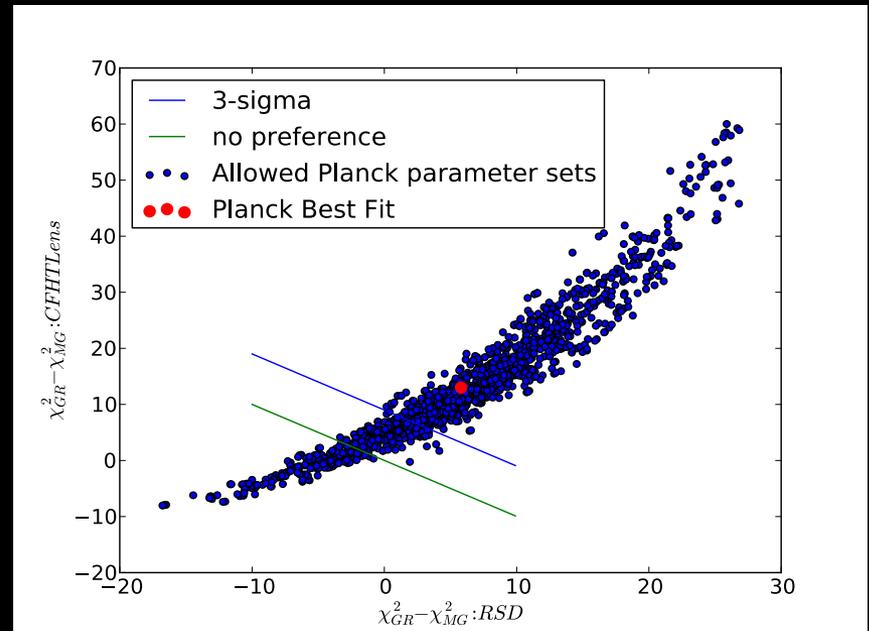
Lensing points are highly correlated, so chi-by-eye fails. Including the proper covariance matrix leads to $\Delta\chi^2=13 \rightarrow$ 3.6-sigma preference for nonlocal model

Allow parameters to vary over Planck range



Allow parameters to vary over Planck range

Naïve interpretation: 68% of allowed region still yield 3-sigma preference; 49% yields 4-sigma, so still very strong preference for nonlocal model

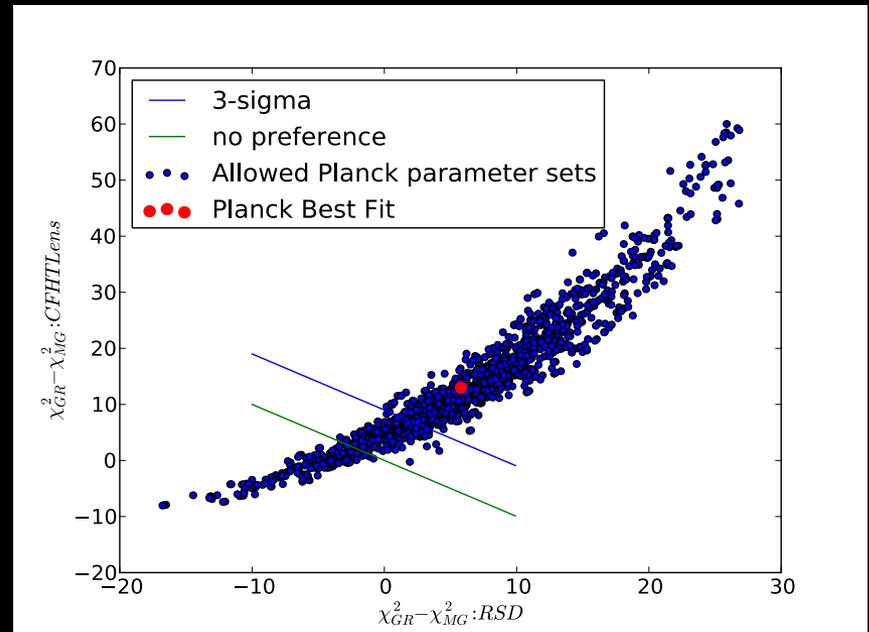


Allow parameters to vary over Planck range

More carefully:

$$\frac{P_{mg} = \int dp P_{mg}(D|p)}{P_{gr} = \int dp P_{gr}(D|p)} = 20 \equiv \exp\left[-\frac{DC^2}{2}\right]$$

2.4-sigma preference for
nonlocal model



Theoretical Challenges in post-Planck era

- Evidence for BSM physics (dark matter, inflation, neutrino masses, dark energy) stronger than ever: *What is this new physics?*
- Is Modified Gravity a viable alternative to dark energy?
- Tension between Planck and galaxy surveys: *Hint that we are not living in a boring universe?*